
empymod Documentation

Release 1.7.3

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The electromagnetic modeller **empymod** can model electric or magnetic responses due to a three-dimensional electric or magnetic source in a layered-earth model with vertical transverse isotropic (VTI) resistivity, VTI electric permittivity, and VTI magnetic permeability, from very low frequencies (DC) to very high frequencies (GPR). The calculation is carried out in the wavenumber-frequency domain, and various Hankel- and Fourier-transform methods are included to transform the responses into the space-frequency and space-time domains.

See <https://empymod.github.io/#features> for a complete list of features.

CHAPTER 1

More information

For more information regarding installation, usage, contributing, roadmap, bug reports, and much more, see

- **Website:** <https://empymod.github.io>,
- **Documentation:** <https://empymod.readthedocs.io>,
- **Source Code:** <https://github.com/empymod>,
- **Examples:** <https://github.com/empymod/example-notebooks>.

CHAPTER 2

Citation

If you publish results for which you used empymod, please give credit by citing this article:

Werthmüller, D., 2017, An open-source full 3D electromagnetic modeler for 1D VTI media in Python: empymod: Geophysics, 82(6), WB9–WB19; DOI: [10.1190/geo2016-0626.1](https://doi.org/10.1190/geo2016-0626.1).

All releases have a Zenodo-DOI, provided on the [release-page](#). Also consider citing [Hunziker et al. \(2015\)](#) and [Key \(2012\)](#), without which empymod would not exist.

License information

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3.1 Manual

3.1.1 Theory

The code is principally based on

- [\[Hunziker_et_al_2015\]](#) for the wavenumber-domain calculation (`kernel`),
- [\[Key_2012\]](#) for the DLF and QWE transforms,
- [\[Slob_et_al_2010\]](#) for the analytical half-space solutions, and
- [\[Hamilton_2000\]](#) for the FFTLog.

See these publications and all the others given in the [references](#), if you are interested in the theory on which `empymod` is based. Another good reference is [\[Ziolkowski_and_Slob_2018\]](#), which will be published in late 2018. The book derives in great detail the equations for layered-Earth CSEM modelling.

3.1.2 Installation

You can install `empymod` either via `conda`:

```
conda install -c prisae empymod
```

or via `pip`:

```
pip install empymod
```

Required are Python version 3.4 or higher and the modules `NumPy` and `SciPy`. The module `numexpr` is required additionally (built with Intel's VML) if you want to run parts of the kernel in parallel.

The modeller `empymod` comes with add-ons (`empymod.scripts`). These add-ons provide some very specific, additional functionalities. Some of these add-ons have additional, optional dependencies for other modules such as `matplotlib`. See the *Add-ons*-section for their documentation.

If you are new to Python I recommend using a Python distribution, which will ensure that all dependencies are met, specifically properly compiled versions of `NumPy` and `SciPy`; I recommend using [Anaconda](https://www.anaconda.com/download). If you install [\[Anaconda\]](https://www.anaconda.com/download)(<https://www.anaconda.com/download>). If you install `Anaconda` you can simply start the *Anaconda Navigator*, add the channel `prisae` and `empymod` will appear in the package list and can be installed with a click.

Warning: Do not use `scipy == 0.19.0`. It has a memory leak in `quad`, see github.com/scipy/scipy/pull/7216. So if you use `QUAD` (or potentially `QWE`) in any of your transforms you might see your memory usage going through the roof.

The structure of `empymod` is:

- **model.py**: EM modelling routines.
- **utils.py**: Utilities for `model` such as checking input parameters.
- **kernel.py**: Kernel of `empymod`, calculates the wavenumber-domain electromagnetic response. Plus analytical, frequency-domain full- and half-space solutions.
- **transform.py**: Methods to carry out the required Hankel transform from wavenumber to space domain and Fourier transform from frequency to time domain.
- **filters.py**: Filters for the *Digital Linear Filters* method DLF (Hankel and Fourier transforms).

3.1.3 Usage/Examples

A good starting point is [\[Werthmuller_2017b\]](#), and more information can be found in [\[Werthmuller_2017\]](#). There are a lot of examples of its usage available, in the form of Jupyter notebooks. Have a look at the following repositories:

- Example notebooks: <https://github.com/empymod/example-notebooks>,
- Geophysical Tutorial TLE: <https://github.com/empymod/article-tle2017>, and
- Numerical examples of [\[Ziolkowski_and_Slob_2018\]](#): <https://github.com/empymod/csem-ziolkowski-and-slob>.

The main modelling routines is `bipole`, which can calculate the electromagnetic frequency- or time-domain field due to arbitrary finite electric or magnetic bipole sources, measured by arbitrary finite electric or magnetic bipole receivers. The model is defined by horizontal resistivity and anisotropy, horizontal and vertical electric permittivities and horizontal and vertical magnetic permeabilities. By default, the electromagnetic response is normalized to source and receiver of 1 m length, and source strength of 1 A.

A simple frequency-domain example, with most of the parameters left at the default value:

```
>>> import numpy as np
>>> from empymod import bipole
>>> # x-directed bipole source: x0, x1, y0, y1, z0, z1
>>> src = [-50, 50, 0, 0, 100, 100]
>>> # x-directed dipole source-array: x, y, z, azimuth, dip
>>> rec = [np.arange(1, 11)*500, np.zeros(10), 200, 0, 0]
>>> # layer boundaries
>>> depth = [0, 300, 1000, 1050]
>>> # layer resistivities
>>> res = [1e20, .3, 1, 50, 1]
>>> # Frequency
>>> freq = 1
>>> # Calculate electric field due to an electric source at 1 Hz.
```

```

>>> # [msrc = mrec = True (default)]
>>> EMfield = bipole(src, rec, depth, res, freq, verb=4)
:: empymod START ::
~
depth      [m] : 0 300 1000 1050
res        [Ohm,m] : 1E+20 0.3 1 50 1
aniso      [-] : 1 1 1 1 1
epermH     [-] : 1 1 1 1 1
epermV     [-] : 1 1 1 1 1
mpermH     [-] : 1 1 1 1 1
mpermV     [-] : 1 1 1 1 1
frequency  [Hz] : 1
Hankel      : DLF (Fast Hankel Transform)
  > Filter      : Key 201 (2009)
  > DLF type    : Standard
Kernel Opt. : None
Loop over   : None (all vectorized)
Source(s)   : 1 bipole(s)
  > intpts      : 1 (as dipole)
  > length      [m] : 100
  > x_c         [m] : 0
  > y_c         [m] : 0
  > z_c         [m] : 100
  > azimuth    [°] : 0
  > dip         [°] : 0
Receiver(s) : 10 dipole(s)
  > x           [m] : 500 - 5000 : 10 [min-max; #]
                  : 500 1000 1500 2000 2500 3000 3500 4000 4500 5000
  > y           [m] : 0 - 0 : 10 [min-max; #]
                  : 0 0 0 0 0 0 0 0 0 0
  > z           [m] : 200
  > azimuth    [°] : 0
  > dip         [°] : 0
Required ab's : 11
~
:: empymod END; runtime = 0:00:00.005536 :: 1 kernel call(s)
~
>>> print (EMfield)
[ 1.68809346e-10 -3.08303130e-10j -8.77189179e-12 -3.76920235e-11j
 -3.46654704e-12 -4.87133683e-12j -3.60159726e-13 -1.12434417e-12j
 1.87807271e-13 -6.21669759e-13j 1.97200208e-13 -4.38210489e-13j
 1.44134842e-13 -3.17505260e-13j 9.92770406e-14 -2.33950871e-13j
 6.75287598e-14 -1.74922886e-13j 4.62724887e-14 -1.32266600e-13j]

```

3.1.4 Contributing

New contributions, bug reports, or any kind of feedback is always welcomed! Have a look at the Roadmap-section to get an idea of things that could be implemented. The best way for interaction is at <https://github.com/empymod>. If you prefer to contact me outside of GitHub use the contact form on my personal website, <https://werthmuller.org>.

To install empymod from source, you can download the latest version from GitHub and either add the path to empymod to your python-path variable, or install it in your python distribution via:

```
python setup.py install
```

Please make sure your code follows the pep8-guidelines by using, for instance, the python module `flake8`, and also that your code is covered with appropriate tests. Just get in touch if you have any doubts.

3.1.5 Tests and benchmarks

The modeller comes with a test suite using `pytest`. If you want to run the tests, just install `pytest` and run it within the `empymod-top-directory`.

```
> pip install pytest coveralls pytest-flake8 pytest-mpl
> # and then
> cd to/the/empymod/folder # Ensure you are in the right directory,
> ls -d */                 # your output should look the same.
docs/ empymod/ tests/
> # pytest will find the tests, which are located in the tests-folder.
> # simply run
> pytest --cov=empymod --flake8 --mpl
```

It should run all tests successfully. Please let me know if not!

Note that installations of `empymod` via `conda` or `pip` do not have the test-suite included. To run the test-suite you must download `empymod` from [GitHub](#).

There is also a benchmark suite using *airspeed velocity*, located in the [empymod/asv-repository](#). The results of my machine can be found in the [empymod/bench](#), its rendered version at [empymod.github.io/asv](#).

3.1.6 Transforms

Included **Hankel transforms**:

- Digital Linear Filters *DLF*
- Quadrature with Extrapolation *QWE*
- Adaptive quadrature *QUAD*

Included **Fourier transforms**:

- Digital Linear Filters *DLF*
- Quadrature with Extrapolation *QWE*
- Logarithmic Fast Fourier Transform *FFTLog*
- Fast Fourier Transform *FFT*

Digital Linear Filters

The module `empymod.filters` comes with many DLFs for the Hankel and the Fourier transform. If you want to export one of these filters to plain ascii files you can use the `tofile`-routine of each filter:

```
>>> import empymod
>>> # Load a filter
>>> filt = empymod.filters.wer_201_2018()
>>> # Save it to pure ascii-files
>>> filt.tofile()
>>> # This will save the following three files:
>>> #     ./filters/wer_201_2018_base.txt
>>> #     ./filters/wer_201_2018_j0.txt
>>> #     ./filters/wer_201_2018_j1.txt
```

Similarly, if you want to use an own filter you can do that as well. The filter base and the filter coefficient have to be stored in separate files:

```
>>> import empymod
>>> # Create an empty filter;
>>> # Name has to be the base of the text files
>>> filt = empymod.filters.DigitalFilter('my-filter')
```

```
>>> # Load the ascii-files
>>> filt.fromfile()
>>> # This will load the following three files:
>>> #     ./filters/my-filter_base.txt
>>> #     ./filters/my-filter_j0.txt
>>> #     ./filters/my-filter_j1.txt
>>> # and store them in filt.base, filt.j0, and filt.j1.
```

The path can be adjusted by providing `tofile` and `fromfile` with a `path`-argument.

FFTLog

FFTLog is the logarithmic analogue to the Fast Fourier Transform FFT originally proposed by [Talman_1978]. The code used by `empymod` was published in Appendix B of [Hamilton_2000] and is publicly available at casa.colorado.edu/~ajsh/FFTLog. From the FFTLog-website:

FFTLog is a set of fortran subroutines that compute the fast Fourier or Hankel (= Fourier-Bessel) transform of a periodic sequence of logarithmically spaced points.

FFTlog can be used for the Hankel as well as for the Fourier Transform, but currently `empymod` uses it only for the Fourier transform. It uses a simplified version of the python implementation of FFTLog, `pyfftlog` (github.com/prisae/pyfftlog).

[Haines_and_Jones_1988] proposed a logarithmic Fourier transform (abbreviated by the authors as LFT) for electromagnetic geophysics, also based on [Talman_1978]. I do not know if Hamilton was aware of the work by Haines and Jones. The two publications share as reference only the original paper by Talman, and both cite a publication of Anderson; Hamilton cites [Anderson_1982], and Haines and Jones cite [Anderson_1979]. Hamilton probably never heard of Haines and Jones, as he works in astronomy, and Haines and Jones was published in the *Geophysical Journal*.

Logarithmic FFTs are not widely used in electromagnetics, as far as I know, probably because of the ease, speed, and generally sufficient precision of the digital filter methods with sine and cosine transforms ([Anderson_1975]). However, comparisons show that FFTLog can be faster and more precise than digital filters, specifically for responses with source and receiver at the interface between air and subsurface. Credit to use FFTLog in electromagnetics goes to David Taylor who, in the mid-2000s, implemented FFTLog into the forward modellers of the company Multi-Transient ElectroMagnetic (MTEM Ltd, later Petroleum Geo-Services PGS). The implementation was driven by land responses, where FFTLog can be much more precise than the filter method for very early times.

Notes on Fourier Transform

The Fourier transform to obtain the space-time domain impulse response from the complex-valued space-frequency response can be calculated by either a cosine transform with the real values, or a sine transform with the imaginary part,

$$\begin{aligned} E(r, t)^{\text{Impulse}} &= \frac{2}{\pi} \int_0^\infty \Re[E(r, \omega)] \cos(\omega t) d\omega, \\ &= -\frac{2}{\pi} \int_0^\infty \Im[E(r, \omega)] \sin(\omega t) d\omega, \end{aligned}$$

see, e.g., [Anderson_1975] or [Key_2012]. Quadrature-with-extrapolation, FFTLog, and obviously the sine/cosine-transform all make use of this split.

To obtain the step-on response the frequency-domain result is first divided by $i\omega$, in the case of the step-off response it is additionally multiplied by -1. The impulse-response is the time-derivative of the step-response,

$$E(r, t)^{\text{Impulse}} = \frac{\partial E(r, t)^{\text{step}}}{\partial t}.$$

Using $\frac{\partial}{\partial t} \Leftrightarrow i\omega$ and going the other way, from impulse to step, leads to the division by $i\omega$. (This only holds because we define in accordance with the causality principle that $E(r, t \leq 0) = 0$).

With the sine/cosine transform (`ft='ffht'/'sin'/'cos'`) you can choose which one you want for the impulse responses. For the switch-on response, however, the sine-transform is enforced, and equally the cosine transform for the switch-off response. This is because these two do not need to know the field at time 0, $E(r, t = 0)$.

The Quadrature-with-extrapolation and FFTLog are hard-coded to use the cosine transform for step-off responses, and the sine transform for impulse and step-on responses. The FFT uses the full complex-valued response at the moment.

For completeness sake, the step-on response is given by

$$E(r, t)^{\text{Step-on}} = -\frac{2}{\pi} \int_0^\infty \Im \left[\frac{E(r, \omega)}{i\omega} \right] \sin(\omega t) d\omega ,$$

and the step-off by

$$E(r, t)^{\text{Step-off}} = -\frac{2}{\pi} \int_0^\infty \Re \left[\frac{E(r, \omega)}{i\omega} \right] \cos(\omega t) d\omega .$$

3.1.7 Note on speed, memory, and accuracy

There is the usual trade-off between speed, memory, and accuracy. Very generally speaking we can say that the *DLF* is faster than *QWE*, but *QWE* is much easier on memory usage. *QWE* allows you to control the accuracy. A standard quadrature in the form of *QUAD* is also provided. *QUAD* is generally orders of magnitudes slower, and more fragile depending on the input arguments. However, it can provide accurate results where *DLF* and *QWE* fail.

The kernel can run in parallel using *numexpr*. This option is activated by setting `opt='parallel'`. It is switched off by default.

I am sure *empymod* could be made much faster with cleverer coding style or with the likes of *cython* or *numba*. Suggestions and contributions are welcomed!

Memory

By default *empymod* will try to carry out the calculation in one go, without looping. If your model has many offsets and many frequencies this can be heavy on memory usage. Even more so if you are calculating time-domain responses for many times. If you are running out of memory, you should use either `loop='off'` or `loop='freq'` to loop over offsets or frequencies, respectively. Use `verb=3` to see how many offsets and how many frequencies are calculated internally.

Depths, Rotation, and Bipole

Depths: Calculation of many source and receiver positions is fastest if they remain at the same depth, as they can be calculated in one kernel-call. If depths do change, one has to loop over them. Note: Sources or receivers placed on a layer interface are considered in the upper layer.

Rotation: Sources and receivers aligned along the principal axes *x*, *y*, and *z* can be calculated in one kernel call. For arbitrary oriented di- or bipoles, 3 kernel calls are required. If source and receiver are arbitrary oriented, 9 (3x3) kernel calls are required.

Bipole: Bipoles increase the calculation time by the amount of integration points used. For a source and a receiver bipole with each 5 integration points you need 25 (5x5) kernel calls. You can calculate it in 1 kernel call if you set both integration points to 1, and therefore calculate the bipole as if they were dipoles at their centre.

Example: For 1 source and 10 receivers, all at the same depth, 1 kernel call is required. If all receivers are at different depths, 10 kernel calls are required. If you make source and receivers bipoles with 5 integration points, 250 kernel calls are required. If you rotate the source arbitrary horizontally, 500 kernel calls are required. If you rotate the receivers too, in the horizontal plane, 1'000 kernel calls are required. If you rotate the receivers also vertically, 1'500 kernel calls are required. If you rotate the source vertically too, 2'250 kernel calls are required.

So your calculation will take 2'250 times longer! No matter how fast the kernel is, this will take a long time. Therefore carefully plan how precise you want to define your source and receiver bipoles.

Table 3.1: Example as a table for comparison: 1 source, 10 receiver (one or many frequencies).

kernel calls	source bipole			receiver bipole			
	intpts	azimuth	dip	intpts	azimuth	dip	diff. z
1	1	0/90	0/90	1	0/90	0/90	1
10	1	0/90	0/90	1	0/90	0/90	10
250	5	0/90	0/90	5	0/90	0/90	10
500	5	arb.	0/90	5	0/90	0/90	10
1000	5	arb.	0/90	5	arb.	0/90	10
1500	5	arb.	0/90	5	arb.	arb.	10
2250	5	arb.	arb.	5	arb.	arb.	10

Parallelisation

If `opt = 'parallel'`, a good dozen of the most time-consuming statements are calculated by using the `numexpr` package (<https://github.com/pydata/numexpr/wiki/Numexpr-Users-Guide>). These statements are all in the kernel-functions `greenfct`, `reflections`, and `fields`, and all involve Γ in one way or another, often calculating square roots or exponentials. As Γ has dimensions (`#frequencies`, `#offsets`, `#layers`, `#lambdas`), it can become fairly big.

The package `numexpr` has to be built with Intel's VML, otherwise it won't be used. You can check if it uses VML with

```
>>> import numexpr
>>> numexpr.use_vml
```

The module `numexpr` uses by default all available cores up to a maximum of 8. You can change this behaviour to a lower or a higher value with the following command (in the example it is changed to 4):

```
>>> import numexpr
>>> numexpr.set_num_threads(4)
```

This parallelisation will make `empymod` faster if you calculate a lot of offsets/frequencies at once, but slower for few offsets/frequencies. Best practice is to check first which one is faster. (You can use the benchmark-notebook in the [empymod/example-notebooks-repository](#).)

Lagged Convolution and Splined Transforms

Both Hankel and Fourier DLF have three options, which can be controlled via the `htarg['pts_per_dec']` and `ftarg['pts_per_dec']` parameters:

- `pts_per_dec=0` : *Standard DLF*;
- `pts_per_dec<0` : *Lagged Convolution DLF*: Spacing defined by filter base, interpolation is carried out in the input domain;
- `pts_per_dec>0` : *Splined DLF*: Spacing defined by `pts_per_dec`, interpolation is carried out in the output domain.

Similarly, interpolation can be used for *QWE* by setting `pts_per_dec` to a value bigger than 0.

The Lagged Convolution and Splined options should be used with caution, as they use interpolation and are therefore less precise than the standard version. However, they can significantly speed up *QWE*, and massively speed up *DLF*. Additionally, the interpolated versions minimizes memory requirements a lot. Speed-up is greater if all source-receiver angles are identical. Note that setting `pts_per_dec` to something else than 0 to calculate only one offset (Hankel) or only one time (Fourier) will be slower than using the standard version. Similarly, the standard version is usually the fastest when using the `parallel` option (`numexpr`).

QWE: Good speed-up is also achieved for *QWE* by setting `maxint` as low as possible. Also, the higher `nquad` is, the higher the speed-up will be.

DLF: Big improvements are achieved for long DLF-filters and for many offsets/frequencies (thousands).

Warning: Keep in mind that setting `pts_per_dec` to something else than 0 uses interpolation, and is therefore not as accurate as the standard version. Use with caution and always compare with the standard version to verify if you can apply interpolation to your problem at hand!

Be aware that *QUAD* (Hankel transform) *always* use the splined version and *always* loops over offsets. The Fourier transforms *FFTlog*, *QWE*, and *FFT* always use interpolation too, either in the frequency or in the time domain. With the *DLF* Fourier transform (sine and cosine transforms) you can choose between no interpolation and interpolation (splined or lagged).

The splined versions of *QWE* check whether the ratio of any two adjacent intervals is above a certain threshold (steep end of the wavenumber or frequency spectrum). If it is, it carries out *QUAD* for this interval instead of *QWE*. The threshold is stored in `diff_quad`, which can be changed within the parameter `htarg` and `ftarg`.

For a graphical explanation of the differences between standard DLF, lagged convolution DLF, and splined DLF for the Hankel and the Fourier transforms see the notebook `7a_DLF-Standard-Lagged-Splined` in the [example-notebooks](#) repository.

Looping

By default, you can calculate many offsets and many frequencies all in one go, vectorized (for the *DLF*), which is the default. The `loop` parameter gives you the possibility to force looping over frequencies or offsets. This parameter can have severe effects on both runtime and memory usage. Play around with this factor to find the fastest version for your problem at hand. It ALWAYS loops over frequencies if `ht = 'QWE' / 'QUAD'` or if `ht = 'FHT'` and `pts_per_dec != 0` (Lagged Convolution or Splined Hankel DLF). All vectorized is very fast if there are few offsets or few frequencies. If there are many offsets and many frequencies, looping over the smaller of the two will be faster. Choosing the right looping together with `opt = 'parallel'` can have a huge influence.

Vertical components

It is advised to use `xdirect = True` (the default) if source and receiver are in the same layer to calculate

- the vertical electric field due to a vertical electric source,
- configurations that involve vertical magnetic components (source or receiver),
- all configurations when source and receiver depth are exactly the same.

The Hankel transforms methods are having sometimes difficulties transforming these functions.

However, in all other cases it will be faster if you set `xdirect = False`.

Time-domain land CSEM

The derivation, as it stands, has a near-singular behaviour in the wavenumber-frequency domain when $\kappa^2 = \omega^2 \epsilon \mu$. This can be a problem for land-domain CSEM calculations if source and receiver are located at the surface between air and subsurface. Because most transforms do not sample the wavenumber-frequency domain sufficiently to catch this near-singular behaviour (hence not smooth), which then creates noise at early times where the signal should be zero. To avoid the issue simply set `epermH[0] = epermV[0] = 0`, hence the relative electric permittivity of the air to zero. This trick obviously uses the diffusive approximation for the air-layer, it therefore will not work for very high frequencies (e.g., GPR calculations).

3.1.8 License

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Note: This software was initially (till 01/2017) developed with funding from *The Mexican National Council of Science and Technology (Consejo Nacional de Ciencia y Tecnología, <http://www.conacyt.gob.mx>)*, carried out at *The Mexican Institute of Petroleum IMP (Instituto Mexicano del Petróleo, <http://www.gob.mx/imp>)*.

3.1.9 References

3.2 Roadmap

A collection of ideas of what could be added or improved in empymod. Please get in touch if you would like to tackle one of these problems!

- **Additional modelling routines**
 - **t_{dem}** (**TEM**) [empymod#8]: Issues that have to be addressed: ramp waveform, windowing, loop integration, zero-offset (coincident loop).
 - * in-loop
 - * coincident loop
 - * loop-loop
 - * arbitrary shaped loops
 - **Ramp waveform** [empymod#7]
 - **Arbitrary waveform** [empymod#7]
 - Improve the GPR-routine [empymod#9]
 - Load and save functions to easily store and load model information (resistivity model, acquisition parameters, and modelling parameters) together with the modelling data (using `pickle` or `shelve`). Probably easier after implementation of the abstraction [empymod#14].
- **Inversion** [empymod#20]: Inversion routines, preferably a selection of different ones.
 - Add some clever checks, e.g. as in Key (2012): abort loops if the field is strongly attenuated.
- Additional (semi-)analytical functions (where possible)
 - Complete full-space (electric and magnetic source and receiver); space-time domain
 - Extend diffusive half-space solution to magnetic sources and receivers; space-frequency and space-time domains
 - Complete half-space
- Transforms
 - Fourier

- * Change `fft` to use discrete sine/cosine transforms instead, as all other Fourier transforms
- * If previous step is successful, clean up the internal decisions (`utils.check_time`) when to use sine/cosine transform (not consistent at the moment, some choice only exists with `ffht` impulse responses, `fqwe` and `fftlog` use sine for impulse, and all three use sine for step-on responses and cosine for step-off responses)
- Hankel
 - * Add the `fht`-module from `FFTLog` for the Hankel transform.
- Hankel and Fourier
 - * Include the method outlined by Mulder et al., 2008, Geophysics (piecewise-cubic Hermite interpolation with a FFT) to try to further speed-up the splined versions.
- Extend examples (example-notebooks)
 - Add different methods (e.g. DC)
 - Reproduce published results
- A `cython`, `numba`, or pure C/C++ implementation of the `kernel` and the `transform` modules. Maybe not worth it, as it may improve speed, but decrease readability. Both at the same time would be nice. A fast C/C++-version for calculations (inversions), and a Python-version to tinker with for interested folks. (Probably combined with default parallelisation, removing the `numexpr` variant.) *Probably not. See closed issue [empymod#21].*
- Abstraction of the code [empymod#14].
- GUI.
- Move `empymod` from channel ‘`prisae`’ to ‘`conda-forge`’ (pros/cons?).

3.3 Changelog

3.3.1 v1.7.3 - 2018-07-16

- Small improvements related to speed as a result of the benchmarks introduced in v1.7.2:
 - Kernels which do not exist for a given `ab` are now returned as `None` from `kernel.wavenumber` instead of arrays of zeroes. This permits for some time saving in the transforms. This change is backwards incompatible if you directly used `kernel.wavenumber`. Nothing changes for the user-facing routines in `model`.
 - Adjustments in `transform` with regard to the `None` returned by `kernel.wavenumber`. The kernels are not checked anymore if they are all zeroes (which can be slow for big arrays). If they are not `None`, they will be processed.
 - Various small improvements for speed to `transform.dlf` (i.e. `factAng`; `log10/log`; re-arranging).

3.3.2 v1.7.2 - 2018-07-07

- Benchmarks: `empymod` has now a benchmark suite, see [empymod/asv](#).
- Fixed a bug in `bipole` for time-domain responses with several receivers or sources with different depths. (Simply failed, as wrong dimension was provided to `tem`).
- Small improvements:
 - Various simplifications or cleaning of the code base.
 - Small change (for speed) in check if kernels are empty in `transform.dlf` and `transform.qwe`.

3.3.3 v1.7.1 - 2018-06-19

- New routines in `empymod.filters.DigitalFilter`: Filters can now be saved to or loaded from pure ascii-files.
- Filters and inversion result from `empymod.scripts.fdesign` are now by default saved in plain text. The filters with their internal routine, the inversion result with `np.savetxt`. Compressed saving can be achieved by giving a name with a '.gz'-ending.
- Change in `empymod.utils`:
 - Renamed `_min_param` to `_min_res`.
 - Anisotropy `aniso` is no longer directly checked for its minimum value. Instead, `res*aniso**2`, hence vertical resistivity, is checked with `_min_res`, and anisotropy is subsequently re-calculated from it.
 - The parameters `epermH`, `epermV`, `mpermH`, and `mpermV` can now be set to 0 (or any positive value) and do not depend on `_min_param`.
- `printinfo`: Generally improved; prints now MKL-info (if available) independently of `numexpr`.
- Simplification of `kernel.reflections` through re-arranging.
- Bug fixes
- Version of re-submission of the DLF article to geophysics.

3.3.4 v1.7.0 - 2018-05-23

Merge `empyscripts` into `empymod` under `empymod.scripts`.

- Clear separation between mandatory and optional imports:
 - Mandatory:
 - * `numpy`
 - * `scipy`
 - Optional:
 - * `numexpr` (for `empymod.kernel`)
 - * `matplotlib` (for `empymod.scripts.fdesign`)
 - * `IPython` (for `empymod.scripts.printinfo`)
- Broaden namespace of `empymod`. All public functions from the various modules and the modules from `empymod.scripts` are now available under `empymod` directly.

3.3.5 v1.6.2 - 2018-05-21

These changes should make calculations using `QWE` and `QUAD` for the Hankel transform for cases which do not require all kernels faster; sometimes as much as twice as fast. However, it might make calculations which do require all kernels a tad slower, as more checks had to be included. (Related to [\[empymod#11\]](#); basically including for `QWE` and `QUAD` what was included for `DLF` in version 1.6.0.)

- `transform`:
 - `dlf`:
 - * Improved by avoiding unnecessary multiplications/summations for empty kernels and applying the angle factor only if it is not 1.
 - * Empty/unused kernels can now be input as `None`, e.g. `signal=(PJ0, None, None)`.
 - * `factAng` is new optional for the Hankel transform, as is `ab`.

- hqwe: Avoids unnecessary calculations for zero kernels, improving speed for these cases.
- hquad, quad: Avoids unnecessary calculations for zero kernels, improving speed for these cases.
- kernel:
 - Simplify wavenumber
 - Simplify angle_factor

3.3.6 v1.6.1 - 2018-05-05

Secondary field calculation.

- Add the possibility to calculate secondary fields only (excluding the direct field) by passing the argument `xdirect=None`. The complete `xdirect`-signature is now (only affects calculation if `src` and `rec` are in the same layer):
 - If `True`, direct field is calculated analytically in the frequency domain.
 - If `False`, direct field is calculated in the wavenumber domain.
 - If `None`, direct field is excluded from the calculation, and only reflected fields are returned (secondary field).
- Bugfix in `model.analytical` for `ab=[36, 63]` (zeroes) [empymod#16].

3.3.7 v1.6.0 - 2018-05-01

This release is not completely backwards compatible for the main modelling routines in `empymod.model`, but almost. Read below to see which functions are affected.

- Improved Hankel DLF [empymod#11]. `empymod.kernel.wavenumber` always returns three kernels, `PJ0`, `PJ1`, and `PJ0b`. The first one is angle-independent, the latter two depend on the angle. Now, depending of what source-receiver configuration is chosen, some of these might be zero. If-statements were now included to avoid the calculation of the DLF, interpolation, and reshaping for 0-kernels, which improves speed for these cases.
- Unified DLF arguments [empymod#10].

These changes are backwards compatible for all main modelling routines in `empymod.model`. However, they are not backwards compatible for the following routines:

- `empymod.model.fem` (removed `use_spline`),
- `empymod.transform.fht` (removed `use_spline`),
- `empymod.transform.hqwe` (removed `use_spline`),
- `empymod.transform.quad` (removed `use_spline`),
- `empymod.transform.dlf` (lagged, splined => `pts_per_dec`),
- `empymod.utils.check_opt` (no longer returns `use_spline`),
- `empymod.utils.check_hankel` (changes in `pts_per_dec`), and
- `empymod.utils.check_time` (changes in `pts_per_dec`).

The function `empymod.utils.spline_backwards_hankel` can be used for backwards compatibility.

Now the Hankel and Fourier DLF have the same behaviour for `pts_per_dec`:

- `pts_per_dec = 0`: Standard DLF,
- `pts_per_dec < 0`: Lagged Convolution DLF, and
- `pts_per_dec > 0`: Splined DLF.

There is one exception which is not backwards compatible: Before, if `opt=None` and `htarg={pts_per_dec: != 0}`, the `pts_per_dec` was not used for the FHT and the QWE. Now, this will be used according to the above definitions.

- Bugfix in `model.wavenumber` for `ab=[36, 63]` (zeroes).

3.3.8 v1.5.2 - 2018-04-25

- DLF improvements:
 - Digital linear filter (DLF) method for the Fourier transform can now be carried out without spline, providing 0 for `pts_per_dec` (or any integer smaller than 1).
 - Combine kernel from `fht` and `ffht` into `dlf`, hence separate DLF from other calculations, as is done with QWE (`qwe` for `hqwe` and `fqwe`).
 - Bug fix regarding `transform.get_spline_values`; a DLF with `pts_per_dec` can now be shorter than the corresponding filter.

3.3.9 v1.5.1 - 2018-02-24

- Documentation:
 - Simplifications: avoid duplication as much as possible between the website (empymod.github.io), the manual (empymod.readthedocs.io), and the README (github.com/empymod/empymod).
 - * Website has now only *Features* and *Installation* in full, all other information comes in the form of links.
 - * README has only information in the form of links.
 - * Manual contains the README, and is basically the main document for all information.
 - Improvements: Change some remaining `md-syntax` to `rst-syntax`.
 - FHT -> DLF: replace FHT as much as possible, without breaking backwards compatibility.

3.3.10 v1.5.0 - 2018-01-02

- Minimum parameter values can now be set and verified with `utils.set_minimum` and `utils.get_minimum`.
- New Hankel filter `wer_201_2018`.
- `opt=parallel` has no effect if `numexpr` is not built against Intel's VML. (Use `import numexpr; numexpr.use_vml` to see if your `numexpr` uses VML.)
- Bug fixes
- Version of manuscript submission to geophysics for the DLF article.

3.3.11 v1.4.4 - 2017-09-18

[This was meant to be 1.4.3, but due to a setup/pypi/anaconda-issue I had to push it to 1.4.4; so there isn't really a version 1.4.3.]

- Add TE/TM split to diffusive ee-halfspace solution.
- Improve `kernel.wavenumber` for fullspaces.

- Extended `fQWE` and `fftlog` to be able to use the cosine-transform. Now the cosine-transform with the real-part frequency response is used internally if a switch-off response (`signal=-1`) is required, rather than calculating the switch-on response (with sine-transform and imaginary-part frequency response) and subtracting it from the DC value.
- Bug fixes
- Version of CSEM book.

3.3.12 v1.4.2 - 2017-06-04

- Bugfix: Fixed squeeze in `model.analytical` with `solution='dsplit'`.
- Version of final submission of manuscript to Geophysics.

3.3.13 v1.4.1 - 2017-05-30

[This was meant to be 1.4.0, but due to a setup/pypi/anaconda-issue I had to push it to 1.4.1; so there isn't really a version 1.4.0.]

- New home: empymod.github.io as entry point, and the project page on github.com/empymod. All `empymod`-repos moved to the new home.
 - `/prisae/empymod` -> `/empymod/empymod`
 - `/prisae/empymod-notebooks` -> `/empymod/example-notebooks`
 - `/prisae/empymod-geo2017` -> `/empymod/article-geo2017`
 - `/prisae/empymod-tle2017` -> `/empymod/article-tle2017`
- Modelling routines:
 - New modelling routine `model.analytical`, which serves as a front-end to `kernel.fullspace` or `kernel.halfspace`.
 - Remove legacy routines `model.time` and `model.frequency`. They are covered perfectly by `model.dipole`.
 - Improved switch-off response (calculate and subtract from DC).
 - `xdirect` adjustments:
 - * `isfullspace` now respects `xdirect`.
 - * Removed `xdirect` from `model.wavenumber` (set to `False`).
- Kernel:
 - Modify `kernel.halfspace` to use same input as other kernel functions.
 - Include time-domain ee halfspace solution into `kernel.halfspace`; possible to obtain direct, reflected, and airwave separately, as well as only fullspace solution (all for the diffusive approximation).

3.3.14 v1.3.0 - 2017-03-30

- Add additional transforms and improve QWE:
 - Conventional adaptive quadrature (QUADPACK) for the Hankel transform;
 - Conventional FFT for the Fourier transform.
 - Add `diff_quad` to `htarg/ftarg` of QWE, a switch parameter for QWE/QUAD.
 - Change QWE/QUAD switch from comparing first interval to comparing all intervals.
 - Add parameters for QUAD (`a`, `b`, `limit`) into `htarg/ftarg` for QWE.

- Allow `htarg/ftarg` as dict additionally to list/tuple.
- Improve `model.gpr`.
- Internal changes:
 - Rename internally the sine/cosine filter from `fft` to `ffht`, because of the addition of the Fast Fourier Transform `fft`.
- Clean-up repository
 - Move `notebooks` to `/prisae/empymod-notebooks`
 - Move `publications/Geophysics2017` to `/prisae/empymod-geo2017`
 - Move `publications/TheLeadingEdge2017` to `/prisae/empymod-tle2017`
- Bug fixes and documentation improvements

3.3.15 v1.2.1 - 2017-03-11

- Change default filter from `key_401_2009` to `key_201_2009` (because of warning regarding 401 pt filter in source code of `DIPOLE1D`.)
- Since 06/02/2017 installable via `pip/conda`.
- Bug fixes

3.3.16 v1.2.0 - 2017-02-02

- New routine:
 - General modelling routine `bipole` (replaces `srcbipole`): Model the EM field for arbitrarily oriented, finite length bipole sources and receivers.
- Added a test suite:
 - Unit-tests of small functions.
 - Framework-tests of the bigger functions:
 - * Comparing to status quo (regression tests),
 - * Comparing to known analytical solutions,
 - * Comparing different options to each other,
 - * Comparing to other 1D modellers (`EMmod`, `DIPOLE1D`, `GREEN3D`).
 - Incorporated with Travis CI and Coveralls.
- Internal changes:
 - Add kernel count (printed if `verb > 1`).
 - `numexpr` is now only required if `opt=='parallel'`. If `numexpr` is not found, `opt` is reset to `None` and a warning is printed.
 - Cleaned-up wavenumber-domain routine.
 - `theta/phi` -> `azimuth/dip`; easier to understand.
 - Refined verbosity levels.
 - Lots of changes in `utils`, with regards to the new routine `bipole` and with regards to verbosity. Moved all warnings out from `transform` and `model` into `utils`.
- Bug fixes

3.3.17 v1.1.0 - 2016-12-22

- New routines:
 - New `srcbipole` modelling routine: Model an arbitrarily oriented, finite length bipole source.
 - Merge `frequency` and `time` into `dipole`. (`frequency` and `time` are still available.)
 - `dipole` now supports multiple sources.
- Internal changes:
 - Replace `get_Gauss_Weights` with `scipy.special.p_roots`
 - `jv(0, x), jv(1, x) -> j0(x), j1(x)`
 - Replace `param_shape` in `utils` with `_check_var` and `_check_shape`.
 - Replace `xco` and `yco` by `angle` in `kernel.fullspace`
 - Replace `fftlog` with python version.
 - Additional sine-/cosine-filters: `key_81_CosSin_2009`, `key_241_CosSin_2009`, and `key_601_CosSin_2009`.
- Bug fixes

3.3.18 v1.0.0 - 2016-11-29

- Initial release; state of manuscript submission to geophysics.

3.4 Credits

Thanks to

- **Jürg Hunziker**, **Kerry Key**, and **Evert Slob** for answering all my questions regarding their codes and publications ([Hunziker et al., 2015](#), [Key, 2009](#), [Key, 2012](#), [Slob et al., 2010](#)).
- **Evert Slob** for the feedback and interaction during the creation of the add-on `tmtmod`, which was developed for the creation of github.com/empymod/csem-ziolkowski-and-slob.
- **Kerry Key** and **Evert Slob** for their inputs and feedback during the development of the add-on `fdesign` (see github.com/empymod/article-fdesign).

3.5 Code

3.5.1 model – Model EM-responses

EM-modelling routines. The implemented routines might not be the fastest solution to your specific problem. Use these routines as template to create your own, problem-specific modelling routine!

Principal routines:

- `bipole`
- `dipole`

The main routine is `bipole`, which can model bipole source(s) and bipole receiver(s) of arbitrary direction, for electric or magnetic sources and receivers, both in frequency and in time. A subset of `bipole` is `dipole`, which models infinitesimal small dipoles along the principal axes `x`, `y`, and `z`.

Further routines are:

- `analytical`: Calculate analytical fullspace and halfspace solutions.

- `wavenumber`: Calculate the electromagnetic wavenumber-domain solution.
- `gpr`: Calculate the Ground-Penetrating Radar (GPR) response.

The `wavenumber` routine can be used if you are interested in the wavenumber-domain result, without Hankel nor Fourier transform. It calls straight the `kernel`. The `gpr`-routine convolves the frequency-domain result with a wavelet, and applies a gain to the time-domain result. This function is still experimental.

The modelling routines make use of the following two core routines:

- **fem**: Calculate wavenumber-domain electromagnetic field and carry out the Hankel transform to the frequency domain.
- **tem**: Carry out the Fourier transform to time domain after `fem`.

`empymod.model.bipole` (*src, rec, depth, res, freqtime, signal=None, aniso=None, epermH=None, epermV=None, mpermH=None, mpermV=None, msrc=False, srcpts=1, mrec=False, recpts=1, strength=0, xdirect=True, ht='fht', htarg=None, ft='sin', ftarg=None, opt=None, loop=None, verb=2*)

Return the electromagnetic field due to an electromagnetic source.

Calculate the electromagnetic frequency- or time-domain field due to arbitrary finite electric or magnetic bipole sources, measured by arbitrary finite electric or magnetic bipole receivers. By default, the electromagnetic response is normalized to to source and receiver of 1 m length, and source strength of 1 A.

Parameters `src, rec` : list of floats or arrays

Source and receiver coordinates (m):

- `[x0, x1, y0, y1, z0, z1]` (bipole of finite length)
- `[x, y, z, azimuth, dip]` (dipole, infinitesimal small)

Dimensions:

- The coordinates `x`, `y`, and `z` (dipole) or `x0`, `x1`, `y0`, `y1`, `z0`, and `z1` (bipole) can be single values or arrays.
- The variables `x` and `y` (dipole) or `x0`, `x1`, `y0`, and `y1` (bipole) must have the same dimensions.
- The variable `z` (dipole) or `z0` and `z1` (bipole) must either be single values or having the same dimension as the other coordinates.
- The variables `azimuth` and `dip` must be single values. If they have different angles, you have to use the bipole-method (with `srcpts/recpts = 1`, so it is calculated as dipoles).

Angles (coordinate system is left-handed, positive `z` down (East-North-Depth):

- `azimuth` (°): horizontal deviation from `x`-axis, anti-clockwise.
- `dip` (°): vertical deviation from `xy`-plane downwards.

Sources or receivers placed on a layer interface are considered in the upper layer.

depth : list

Absolute layer interfaces `z` (m); `#depth = #res - 1` (excluding +/- infinity).

res : array_like

Horizontal resistivities `rho_h` (Ohm.m); `#res = #depth + 1`.

freqtime : array_like

Frequencies `f` (Hz) if `signal == None`, else times `t` (s); (`f`, `t` > 0).

signal : {None, 0, 1, -1}, optional

Source signal, default is None:

- None: Frequency-domain response

- -1 : Switch-off time-domain response
- 0 : Impulse time-domain response
- +1 : Switch-on time-domain response

aniso : array_like, optional

Anisotropies $\lambda = \sqrt{\rho_v/\rho_h}$ (-); #aniso = #res. Defaults to ones.

epermH, epermV : array_like, optional

Relative horizontal/vertical electric permittivities ϵ_h/ϵ_v (-); #epermH = #epermV = #res. Default is ones.

mpermH, mpermV : array_like, optional

Relative horizontal/vertical magnetic permeabilities μ_h/μ_v (-); #mpermH = #mpermV = #res. Default is ones.

msrc, mrec : boolean, optional

If True, source/receiver (msrc/mrec) is magnetic, else electric. Default is False.

srcpts, recpts : int, optional

Number of integration points for bipole source/receiver, default is 1:

- srcpts/recpts < 3 : bipole, but calculated as dipole at centre
- srcpts/recpts >= 3 : bipole

strength : float, optional

Source strength (A):

- If 0, output is normalized to source and receiver of 1 m length, and source strength of 1 A.
- If != 0, output is returned for given source and receiver length, and source strength.

Default is 0.

xdirect : bool or None, optional

Direct field calculation (only if src and rec are in the same layer):

- If True, direct field is calculated analytically in the frequency domain.
- If False, direct field is calculated in the wavenumber domain.
- If None, direct field is excluded from the calculation, and only reflected fields are returned (secondary field).

Defaults to True.

ht : { 'fht', 'qwe', 'quad' }, optional

Flag to choose either the *Digital Linear Filter* method (FHT, *Fast Hankel Transform*), the *Quadrature-With-Extrapolation* (QWE), or a simple *Quadrature* (QUAD) for the Hankel transform. Defaults to 'fht'.

htarg : dict or list, optional

Depends on the value for ht:

- If ht = 'fht': [fhtfilt, pts_per_dec]:
 - **fhtfilt**: string of filter name in `empymod.filters` or the filter method itself. (default: `empymod.filters.key_201_2009()`)
 - **pts_per_dec**: points per decade; (default: 0)
 - * If 0: Standard DLF.

- * If < 0 : Lagged Convolution DLF.
- * If > 0 : Splined DLF
- If `ht = 'qwe'`: [`rtol`, `atol`, `nquad`, `maxint`, `pts_per_dec`, `diff_quad`, `a`, `b`, `limit`]:
 - `rtol`: relative tolerance (default: $1e-12$)
 - `atol`: absolute tolerance (default: $1e-30$)
 - `nquad`: order of Gaussian quadrature (default: 51)
 - **`maxint`: maximum number of partial integral intervals** (default: 40)
 - **`pts_per_dec`: points per decade; (default: 0)**
 - * If 0, no interpolation is used.
 - * If > 0 , interpolation is used.
 - `diff_quad`: criteria when to swap to QUAD (only relevant if `opt='spline'`) (default: 100)
 - `a`: lower limit for QUAD (default: first interval from QWE)
 - `b`: upper limit for QUAD (default: last interval from QWE)
 - `limit`: limit for quad (default: `maxint`)
- If `ht = 'quad'`: [`atol`, `rtol`, `limit`, `lmin`, `lmax`, `pts_per_dec`]:
 - `rtol`: relative tolerance (default: $1e-12$)
 - `atol`: absolute tolerance (default: $1e-20$)
 - `limit`: An upper bound on the number of subintervals used in the adaptive algorithm (default: 500)
 - `lmin`: Minimum wavenumber (default $1e-6$)
 - `lmax`: Maximum wavenumber (default 0.1)
 - `pts_per_dec`: points per decade (default: 40)

The values can be provided as dict with the keywords, or as list. However, if provided as list, you have to follow the order given above. A few examples, assuming `ht = qwe`:

- **Only changing `rtol`**: {`'rtol': 1e-4`} or [`1e-4`] or `1e-4`
- **Changing `rtol` and `nquad`**: {`'rtol': 1e-4`, `'nquad': 101`} or [`1e-4`, `'`, 101]
- **Only changing `diff_quad`**: {`'diffquad': 10`} or [`'`, `'`, `'`, `'`, `'`, 10]

`ft`: {`'sin'`, `'cos'`, `'qwe'`, `'fftlog'`, `'fft'`}, optional

Only used if `signal != None`. Flag to choose either the Digital Linear Filter method (Sine- or Cosine-Filter), the Quadrature-With-Extrapolation (QWE), the FFTLog, or the FFT for the Fourier transform. Defaults to `'sin'`.

`ftarg`: dict or list, optional

Only used if `signal != None`. Depends on the value for `ft`:

- If `ft = 'sin'` or `'cos'`: [`fftfilt`, `pts_per_dec`]:
 - **`fftfilt`: string of filter name in `empymod.filters` or the filter method itself.** (Default: `empymod.filters.key_201_CosSin_2012()`)

- **pts_per_dec: points per decade; (default: -1)**
 - * If 0: Standard DLF.
 - * If < 0: Lagged Convolution DLF.
 - * If > 0: Splined DLF
- If `fft = 'qwe'`: [rtol, atol, nquad, maxint, pts_per_dec]:
 - rtol: relative tolerance (default: 1e-8)
 - atol: absolute tolerance (default: 1e-20)
 - nquad: order of Gaussian quadrature (default: 21)
 - **maxint: maximum number of partial integral intervals** (default: 200)
 - pts_per_dec: points per decade (default: 20)
 - diff_quad: criteria when to swap to QUAD (default: 100)
 - a: lower limit for QUAD (default: first interval from QWE)
 - b: upper limit for QUAD (default: last interval from QWE)
 - limit: limit for quad (default: maxint)
- If `fft = 'fftlog'`: [pts_per_dec, add_dec, q]:
 - pts_per_dec: sampels per decade (default: 10)
 - add_dec: additional decades [left, right] (default: [-2, 1])
 - q: exponent of power law bias (default: 0); -1 <= q <= 1
- If `fft = 'fft'`: [dfreq, nfreq, ntot]:
 - dfreq: Linear step-size of frequencies (default: 0.002)
 - nfreq: Number of frequencies (default: 2048)
 - **ntot: Total number for FFT; difference between nfreq and ntot is padded with zeroes.** This number is ideally a power of 2, e.g. 2048 or 4096 (default: nfreq).
 - pts_per_dec : points per decade (default: None)

Padding can sometimes improve the result, not always. The default samples from 0.002 Hz - 4.096 Hz. If `pts_per_dec` is set to an integer, calculated frequencies are logarithmically spaced with the given number per decade, and then interpolated to yield the required frequencies for the FFT.

The values can be provided as dict with the keywords, or as list. However, if provided as list, you have to follow the order given above. See `htarg` for a few examples.

opt : {None, 'parallel'}, optional

Optimization flag. Defaults to None:

- None: Normal case, no parallelization nor interpolation is used.
- If 'parallel', the package `numexpr` is used to evaluate the most expensive statements. Always check if it actually improves performance for a specific problem. It can speed up the calculation for big arrays, but will most likely be slower for small arrays. It will use all available cores for these specific statements, which all contain `Gamma` in one way or another, which has dimensions (`#frequencies`, `#offsets`, `#layers`, `#lambdas`), therefore can grow pretty big. The module `numexpr`

uses by default all available cores up to a maximum of 8. You can change this behaviour to your desired number of threads `nthreads` with `numexpr.set_num_threads(nthreads)`.

- The value ‘spline’ is deprecated and will be removed. See `htarg` instead for the interpolated versions.

The option ‘parallel’ only affects speed and memory usage, whereas ‘spline’ also affects precision! Please read the note in the *README* documentation for more information.

loop : {None, ‘freq’, ‘off’}, optional

Define if to calculate everything vectorized or if to loop over frequencies (‘freq’) or over offsets (‘off’), default is None. It always loops over frequencies if `ht = 'qwe'` or if `opt = 'spline'`. Calculating everything vectorized is fast for few offsets OR for few frequencies. However, if you calculate many frequencies for many offsets, it might be faster to loop over frequencies. Only comparing the different versions will yield the answer for your specific problem at hand!

verb : {0, 1, 2, 3, 4}, optional

Level of verbosity, default is 2:

- 0: Print nothing.
- 1: Print warnings.
- 2: Print additional runtime and kernel calls
- 3: Print additional start/stop, condensed parameter information.
- 4: Print additional full parameter information

Returns **EM** : ndarray, (nfreq, nrec, nsrc)

Frequency- or time-domain EM field (depending on signal):

- If rec is electric, returns E [V/m].
- If rec is magnetic, returns B [T] (not H [A/m]!).

In the case of the impulse time-domain response, the unit is further divided by seconds [1/s].

However, source and receiver are normalised (unless strength != 0). So for instance in the electric case the source strength is 1 A and its length is 1 m. So the electric field could also be written as [V/(A.m2)].

In the magnetic case the source strength is given by $i\omega\mu_0 AI^e$, where A is the loop area (m2), and I^e the electric source strength. For the normalized magnetic source $A = 1\text{m}^2$ and $I^e = 1\text{Ampere}$. A magnetic source is therefore frequency dependent.

The shape of EM is (nfreq, nrec, nsrc). However, single dimensions are removed.

See also:

fem Electromagnetic frequency-domain response.

tem Electromagnetic time-domain response.

Examples

```
>>> import numpy as np
>>> from empymod import bipole
>>> # x-directed bipole source: x0, x1, y0, y1, z0, z1
>>> src = [-50, 50, 0, 0, 100, 100]
```

```

>>> # x-directed dipole source-array: x, y, z, azimuth, dip
>>> rec = [np.arange(1, 11)*500, np.zeros(10), 200, 0, 0]
>>> # layer boundaries
>>> depth = [0, 300, 1000, 1050]
>>> # layer resistivities
>>> res = [1e20, .3, 1, 50, 1]
>>> # Frequency
>>> freq = 1
>>> # Calculate electric field due to an electric source at 1 Hz.
>>> # [msrc = mrec = True (default)]
>>> EMfield = bipole(src, rec, depth, res, freq, verb=4)
:: empymod START ::
~
depth      [m] : 0 300 1000 1050
res        [Ohm.m] : 1E+20 0.3 1 50 1
aniso      [-] : 1 1 1 1 1
epermH     [-] : 1 1 1 1 1
epermV     [-] : 1 1 1 1 1
mpermH     [-] : 1 1 1 1 1
mpermV     [-] : 1 1 1 1 1
frequency  [Hz] : 1
Hankel      : DLF (Fast Hankel Transform)
  > Filter   : Key 201 (2009)
  > DLF type : Standard
Kernel Opt. : None
Loop over   : None (all vectorized)
Source(s)   : 1 bipole(s)
  > intpts   : 1 (as dipole)
  > length   [m] : 100
  > x_c      [m] : 0
  > y_c      [m] : 0
  > z_c      [m] : 100
  > azimuth  [°] : 0
  > dip      [°] : 0
Receiver(s) : 10 dipole(s)
  > x        [m] : 500 - 5000 : 10 [min-max; #]
                : 500 1000 1500 2000 2500 3000 3500 4000 4500 5000
  > y        [m] : 0 - 0 : 10 [min-max; #]
                : 0 0 0 0 0 0 0 0 0 0
  > z        [m] : 200
  > azimuth  [°] : 0
  > dip      [°] : 0
Required ab's : 11
~
:: empymod END; runtime = 0:00:00.005536 :: 1 kernel call(s)
~
>>> print(EMfield)
[ 1.68809346e-10 -3.08303130e-10j -8.77189179e-12 -3.76920235e-11j
 -3.46654704e-12 -4.87133683e-12j -3.60159726e-13 -1.12434417e-12j
  1.87807271e-13 -6.21669759e-13j  1.97200208e-13 -4.38210489e-13j
  1.44134842e-13 -3.17505260e-13j  9.92770406e-14 -2.33950871e-13j
  6.75287598e-14 -1.74922886e-13j  4.62724887e-14 -1.32266600e-13j]

```

`empymod.model.dipole(src, rec, depth, res, freqtime, signal=None, ab=11, aniso=None, epermH=None, epermV=None, mpermH=None, mpermV=None, xdirect=True, ht='fht', htarg=None, ft='sin', ftarg=None, opt=None, loop=None, verb=2)`

Return the electromagnetic field due to a dipole source.

Calculate the electromagnetic frequency- or time-domain field due to infinitesimal small electric or magnetic dipole source(s), measured by infinitesimal small electric or magnetic dipole receiver(s); sources and receivers are directed along the principal directions x, y, or z, and all sources are at the same depth, as well

as all receivers are at the same depth.

Use the functions `bipole` to calculate dipoles with arbitrary angles or bipoles of finite length and arbitrary angle.

The function `dipole` could be replaced by `bipole` (all there is to do is translate `ab` into `msrc`, `mrec`, `azimuth`'s and `dip`'s). However, `dipole` is kept separately to serve as an example of a simple modelling routine that can serve as a template.

Parameters `src`, `rec` : list of floats or arrays

Source and receiver coordinates (m): [x, y, z]. The x- and y-coordinates can be arrays, z is a single value. The x- and y-coordinates must have the same dimension.

Sources or receivers placed on a layer interface are considered in the upper layer.

depth : list

Absolute layer interfaces z (m); `#depth = #res - 1` (excluding +/- infinity).

res : array_like

Horizontal resistivities `rho_h` (Ohm.m); `#res = #depth + 1`.

freqtime : array_like

Frequencies `f` (Hz) if `signal == None`, else times `t` (s); (`f`, `t` > 0).

signal : {None, 0, 1, -1}, optional

Source signal, default is None:

- None: Frequency-domain response
- -1 : Switch-off time-domain response
- 0 : Impulse time-domain response
- +1 : Switch-on time-domain response

ab : int, optional

Source-receiver configuration, defaults to 11.

		electric source			magnetic source		
		x	y	z	x	y	z
electric receiver	x	11	12	13	14	15	16
	y	21	22	23	24	25	26
	z	31	32	33	34	35	36
magnetic receiver	x	41	42	43	44	45	46
	y	51	52	53	54	55	56
	z	61	62	63	64	65	66

aniso : array_like, optional

Anisotropies `lambda = sqrt(rho_v/rho_h)` (-); `#aniso = #res`. Defaults to ones.

epermH, **epermV** : array_like, optional

Relative horizontal/vertical electric permittivities `epsilon_h/epsilon_v` (-); `#epermH = #epermV = #res`. Default is ones.

mpermH, **mpermV** : array_like, optional

Relative horizontal/vertical magnetic permeabilities `mu_h/mu_v` (-); `#mpermH = #mpermV = #res`. Default is ones.

xdirect : bool or None, optional

Direct field calculation (only if src and rec are in the same layer):

- If True, direct field is calculated analytically in the frequency domain.
- If False, direct field is calculated in the wavenumber domain.
- If None, direct field is excluded from the calculation, and only reflected fields are returned (secondary field).

Defaults to True.

ht : { 'fht', 'qwe', 'quad' }, optional

Flag to choose either the *Digital Linear Filter* method (FHT, *Fast Hankel Transform*), the *Quadrature-With-Extrapolation* (QWE), or a simple *Quadrature* (QUAD) for the Hankel transform. Defaults to 'fht'.

htarg : dict or list, optional

Depends on the value for ht:

- If **ht** = 'fht': [fhtfilt, pts_per_dec]:
 - **fhtfilt**: string of filter name in `empymod.filters` or the filter method itself. (default: `empymod.filters.key_201_2009()`)
 - **pts_per_dec**: points per decade; (default: 0)
 - * If 0: Standard DLF.
 - * If < 0: Lagged Convolution DLF.
 - * If > 0: Splined DLF
- If **ht** = 'qwe': [rtol, atol, nquad, maxint, pts_per_dec, diff_quad, a, b, limit]:
 - **rtol**: relative tolerance (default: 1e-12)
 - **atol**: absolute tolerance (default: 1e-30)
 - **nquad**: order of Gaussian quadrature (default: 51)
 - **maxint**: maximum number of partial integral intervals (default: 40)
 - **pts_per_dec**: points per decade; (default: 0)
 - * If 0, no interpolation is used.
 - * If > 0, interpolation is used.
 - **diff_quad**: criteria when to swap to QUAD (only relevant if opt='spline') (default: 100)
 - **a**: lower limit for QUAD (default: first interval from QWE)
 - **b**: upper limit for QUAD (default: last interval from QWE)
 - **limit**: limit for quad (default: maxint)
- If **ht** = 'quad': [atol, rtol, limit, lmin, lmax, pts_per_dec]:
 - **rtol**: relative tolerance (default: 1e-12)
 - **atol**: absolute tolerance (default: 1e-20)
 - **limit**: An upper bound on the number of subintervals used in the adaptive algorithm (default: 500)
 - **lmin**: Minimum wavenumber (default 1e-6)
 - **lmax**: Maximum wavenumber (default 0.1)
 - **pts_per_dec**: points per decade (default: 40)

The values can be provided as dict with the keywords, or as list. However, if provided as list, you have to follow the order given above. A few examples, assuming **ht** = 'qwe':

- **Only changing rtol**: { 'rtol': 1e-4 } or [1e-4] or 1e-4

- **Changing rtol and nquad:** {'rtol': 1e-4, 'nquad': 101} or [1e-4, '', 101]

- **Only changing diff_quad:** {'diffquad': 10} or ['', '', '', '', '', 10]

ft : {'sin', 'cos', 'qwe', 'fftlog', 'fft'}, optional

Only used if `signal != None`. Flag to choose either the Digital Linear Filter method (Sine- or Cosine-Filter), the Quadrature-With-Extrapolation (QWE), the FFTLog, or the FFT for the Fourier transform. Defaults to 'sin'.

ftarg : dict or list, optional

Only used if `signal != None`. Depends on the value for **ft**:

- If `ft = 'sin' or 'cos'`: [fftfilt, pts_per_dec]:
 - **fftfilt**: string of filter name in `empymod.filters` or the filter method itself. (Default: `empymod.filters.key_201_CosSin_2012()`)
 - **pts_per_dec**: points per decade; (default: -1)
 - * If 0: Standard DLF.
 - * If < 0: Lagged Convolution DLF.
 - * If > 0: Splined DLF
- If `ft = 'qwe'`: [rtol, atol, nquad, maxint, pts_per_dec]:
 - **rtol**: relative tolerance (default: 1e-8)
 - **atol**: absolute tolerance (default: 1e-20)
 - **nquad**: order of Gaussian quadrature (default: 21)
 - **maxint**: maximum number of partial integral intervals (default: 200)
 - **pts_per_dec**: points per decade (default: 20)
 - **diff_quad**: criteria when to swap to QUAD (default: 100)
 - **a**: lower limit for QUAD (default: first interval from QWE)
 - **b**: upper limit for QUAD (default: last interval from QWE)
 - **limit**: limit for quad (default: maxint)
- If `ft = 'fftlog'`: [pts_per_dec, add_dec, q]:
 - **pts_per_dec**: sampels per decade (default: 10)
 - **add_dec**: additional decades [left, right] (default: [-2, 1])
 - **q**: exponent of power law bias (default: 0); -1 <= q <= 1
- If `ft = 'fft'`: [dfreq, nfreq, ntot]:
 - **dfreq**: Linear step-size of frequencies (default: 0.002)
 - **nfreq**: Number of frequencies (default: 2048)
 - **ntot**: Total number for FFT; difference between **nfreq** and **ntot** is padded with zeroes. This number is ideally a power of 2, e.g. 2048 or 4096 (default: **nfreq**).
 - **pts_per_dec** : points per decade (default: None)

Padding can sometimes improve the result, not always. The default samples from 0.002 Hz - 4.096 Hz. If `pts_per_dec` is set to an integer, calculated frequencies are logarithmically spaced with the given number per

decade, and then interpolated to yield the required frequencies for the FFT.

The values can be provided as dict with the keywords, or as list. However, if provided as list, you have to follow the order given above. See `htarg` for a few examples.

opt : {None, 'parallel'}, optional

Optimization flag. Defaults to None:

- None: Normal case, no parallelization nor interpolation is used.
- If 'parallel', the package `numexpr` is used to evaluate the most expensive statements. Always check if it actually improves performance for a specific problem. It can speed up the calculation for big arrays, but will most likely be slower for small arrays. It will use all available cores for these specific statements, which all contain `Gamma` in one way or another, which has dimensions (`#frequencies`, `#offsets`, `#layers`, `#lambdas`), therefore can grow pretty big. The module `numexpr` uses by default all available cores up to a maximum of 8. You can change this behaviour to your desired number of threads `nthreads` with `numexpr.set_num_threads(nthreads)`.
- The value 'spline' is deprecated and will be removed. See `htarg` instead for the interpolated versions.

The option 'parallel' only affects speed and memory usage, whereas 'spline' also affects precision! Please read the note in the *README* documentation for more information.

loop : {None, 'freq', 'off'}, optional

Define if to calculate everything vectorized or if to loop over frequencies ('freq') or over offsets ('off'), default is None. It always loops over frequencies if `ht = 'qwe'` or if `opt = 'spline'`. Calculating everything vectorized is fast for few offsets OR for few frequencies. However, if you calculate many frequencies for many offsets, it might be faster to loop over frequencies. Only comparing the different versions will yield the answer for your specific problem at hand!

verb : {0, 1, 2, 3, 4}, optional

Level of verbosity, default is 2:

- 0: Print nothing.
- 1: Print warnings.
- 2: Print additional runtime and kernel calls
- 3: Print additional start/stop, condensed parameter information.
- 4: Print additional full parameter information

Returns **EM** : ndarray, (nfreq, nrec, nsrc)

Frequency- or time-domain EM field (depending on signal):

- If `rec` is electric, returns `E` [V/m].
- If `rec` is magnetic, returns `B` [T] (not `H` [A/m!]).

In the case of the impulse time-domain response, the unit is further divided by seconds [1/s].

However, source and receiver are normalised. So for instance in the electric case the source strength is 1 A and its length is 1 m. So the electric field could also be written as [V/(A.m²)].

The shape of `EM` is (nfreq, nrec, nsrc). However, single dimensions are removed.

See also:

bipole Electromagnetic field due to an electromagnetic source.

fem Electromagnetic frequency-domain response.

tem Electromagnetic time-domain response.

Examples

```
>>> import numpy as np
>>> from empymod import dipole
>>> src = [0, 0, 100]
>>> rec = [np.arange(1, 11)*500, np.zeros(10), 200]
>>> depth = [0, 300, 1000, 1050]
>>> res = [1e20, .3, 1, 50, 1]
>>> EMfield = dipole(src, rec, depth, res, freqtime=1, verb=0)
>>> print(EMfield)
[ 1.68809346e-10 -3.08303130e-10j -8.77189179e-12 -3.76920235e-11j
 -3.46654704e-12 -4.87133683e-12j -3.60159726e-13 -1.12434417e-12j
 1.87807271e-13 -6.21669759e-13j 1.97200208e-13 -4.38210489e-13j
 1.44134842e-13 -3.17505260e-13j 9.92770406e-14 -2.33950871e-13j
 6.75287598e-14 -1.74922886e-13j 4.62724887e-14 -1.32266600e-13j]
```

```
empymod.model.analytical(src, rec, res, freqtime, solution='fs', signal=None, ab=11,
                          aniso=None, epermH=None, epermV=None, mpermH=None,
                          mpermV=None, verb=2)
```

Return the analytical full- or half-space solution.

Calculate the electromagnetic frequency- or time-domain field due to infinitesimal small electric or magnetic dipole source(s), measured by infinitesimal small electric or magnetic dipole receiver(s); sources and receivers are directed along the principal directions x, y, or z, and all sources are at the same depth, as well as all receivers are at the same depth.

In the case of a halfspace the air-interface is located at $z = 0$ m.

You can call the functions `fullspace` and `halfspace` in `kernel.py` directly. This interface is just to provide a consistent interface with the same input parameters as for instance for `dipole`.

This function yields the same result if `solution='fs'` as `dipole`, if the model is a fullspace.

Included are:

- Full fullspace solution (`solution='fs'`) for ee-, me-, em-, mm-fields, only frequency domain, [Hunziker_et_al_2015].
- Diffusive fullspace solution (`solution='dfs'`) for ee-fields, [Slob_et_al_2010].
- Diffusive halfspace solution (`solution='dhs'`) for ee-fields, [Slob_et_al_2010].
- Diffusive direct- and reflected field and airwave (`solution='dsplit'`) for ee-fields, [Slob_et_al_2010].
- Diffusive direct- and reflected field and airwave (`solution='dtetm'`) for ee-fields, split into TE and TM mode [Slob_et_al_2010].

Parameters `src, rec` : list of floats or arrays

Source and receiver coordinates (m): [x, y, z]. The x- and y-coordinates can be arrays, z is a single value. The x- and y-coordinates must have the same dimension.

`res` : float

Horizontal resistivity ρ_h (Ohm.m).

`freqtime` : array_like

Frequencies f (Hz) if `signal == None`, else times t (s); ($f, t > 0$).

`solution` : str, optional

Defines which solution is returned:

- 'fs' : Full fullspace solution (ee-, me-, em-, mm-fields); f-domain.
- 'dfs' : Diffusive fullspace solution (ee-fields only).

- ‘dhs’ : Diffusive halfspace solution (ee-fields only).
- ‘dsplit’ [Diffusive direct- and reflected field and airwave] (ee-fields only).
- ‘dtetm’ [as dsplit, but direct field TE, TM; reflected field TE, TM,] and airwave (ee-fields only).

signal : {None, 0, 1, -1}, optional

Source signal, default is None:

- None: Frequency-domain response
- -1 : Switch-off time-domain response
- 0 : Impulse time-domain response
- +1 : Switch-on time-domain response

ab : int, optional

Source-receiver configuration, defaults to 11.

		electric source			magnetic source		
		x	y	z	x	y	z
electric receiver	x	11	12	13	14	15	16
	y	21	22	23	24	25	26
	z	31	32	33	34	35	36
magnetic receiver	x	41	42	43	44	45	46
	y	51	52	53	54	55	56
	z	61	62	63	64	65	66

aniso : float, optional

Anisotropy $\lambda = \sqrt{\rho_v/\rho_h}$ (-); defaults to one.

epermH, epermV : float, optional

Relative horizontal/vertical electric permittivity ϵ_h/ϵ_v (-); default is one. Ignored for the diffusive solution.

mpermH, mpermV : float, optional

Relative horizontal/vertical magnetic permeability μ_h/μ_v (-); default is one. Ignored for the diffusive solution.

verb : {0, 1, 2, 3, 4}, optional

Level of verbosity, default is 2:

- 0: Print nothing.
- 1: Print warnings.
- 2: Print additional runtime
- 3: Print additional start/stop, condensed parameter information.
- 4: Print additional full parameter information

Returns **EM** : ndarray, (nfreq, nrec, nsrc)

Frequency- or time-domain EM field (depending on signal):

- If rec is electric, returns E [V/m].
- If rec is magnetic, returns B [T] (not H [A/m]!).

In the case of the impulse time-domain response, the unit is further divided by seconds [1/s].

However, source and receiver are normalised. So for instance in the electric case the source strength is 1 A and its length is 1 m. So the electric field could also be written as [V/(A.m²)].

The shape of EM is (nfreq, nrec, nsrc). However, single dimensions are removed.

If `solution='dsplit'`, three ndarrays are returned: `direct`, `reflect`, `air`.

If `solution='dtetm'`, five ndarrays are returned: `direct_TE`, `direct_TM`, `reflect_TE`, `reflect_TM`, `air`.

Examples

```
>>> import numpy as np
>>> from empymod import analytical
>>> src = [0, 0, 0]
>>> rec = [np.arange(1, 11)*500, np.zeros(10), 200]
>>> res = 50
>>> EMfield = analytical(src, rec, res, freqtime=1, verb=0)
>>> print(EMfield)
[ 4.03091405e-08 -9.69163818e-10j  6.97630362e-09 -4.88342150e-10j
 2.15205979e-09 -2.97489809e-10j  8.90394459e-10 -1.99313433e-10j
 4.32915802e-10 -1.40741644e-10j  2.31674165e-10 -1.02579391e-10j
 1.31469130e-10 -7.62770461e-11j  7.72342470e-11 -5.74534125e-11j
 4.61480481e-11 -4.36275540e-11j  2.76174038e-11 -3.32860932e-11j]
```

```
empymod.model.gpr(src, rec, depth, res, freqtime, cf, gain=None, ab=11, aniso=None,
                  epermH=None, epermV=None, mpermH=None, mpermV=None, xdirect=True,
                  ht='quad', htarg=None, ft='fft', ftarg=None, opt=None, loop=None, verb=2)
```

Return the Ground-Penetrating Radar signal.

THIS FUNCTION IS EXPERIMENTAL, USE WITH CAUTION.

It is rather an example how you can calculate GPR responses; however, DO NOT RELY ON IT! It works only well with QUAD or QWE (`quad`, `qwe`) for the Hankel transform, and with FFT (`fft`) for the Fourier transform.

It calls internally `dipole` for the frequency-domain calculation. It subsequently convolves the response with a Ricker wavelet with central frequency `cf`. If `signal!=None`, it carries out the Fourier transform and applies a gain to the response.

For input parameters see the function `dipole`, except for:

Parameters `cf`: float

Centre frequency of GPR-signal, in Hz. Sensible values are between 10 MHz and 3000 MHz.

gain: float

Power of gain function. If `None`, no gain is applied. Only used if `signal!=None`.

Returns `EM`: ndarray

GPR response

```
empymod.model.wavenumber(src, rec, depth, res, freq, wavenumber, ab=11, aniso=None,
                          epermH=None, epermV=None, mpermH=None, mpermV=None,
                          verb=2)
```

Return the electromagnetic wavenumber-domain field.

Calculate the electromagnetic wavenumber-domain field due to infinitesimal small electric or magnetic dipole source(s), measured by infinitesimal small electric or magnetic dipole receiver(s); sources and receivers are directed along the principal directions x, y, or z, and all sources are at the same depth, as well as all receivers are at the same depth.

Parameters `src, rec`: list of floats or arrays

Source and receiver coordinates (m): [x, y, z]. The x- and y-coordinates can be arrays, z is a single value. The x- and y-coordinates must have the same dimension. The x- and y-coordinates only matter for the angle-dependent factor.

Sources or receivers placed on a layer interface are considered in the upper layer.

depth : list

Absolute layer interfaces z (m); $\#depth = \#res - 1$ (excluding +/- infinity).

res : array_like

Horizontal resistivities ρ_h (Ohm.m); $\#res = \#depth + 1$.

freq : array_like

Frequencies f (Hz), used to calculate $\eta H/V$ and $\zeta H/V$.

wavenumber : array

Wavenumbers λ (1/m)

ab : int, optional

Source-receiver configuration, defaults to 11.

		electric source			magnetic source		
		x	y	z	x	y	z
electric receiver	x	11	12	13	14	15	16
	y	21	22	23	24	25	26
	z	31	32	33	34	35	36
magnetic receiver	x	41	42	43	44	45	46
	y	51	52	53	54	55	56
	z	61	62	63	64	65	66

aniso : array_like, optional

Anisotropies $\lambda = \sqrt{\rho_v/\rho_h}$ (-); $\#aniso = \#res$. Defaults to ones.

epermH, epermV : array_like, optional

Relative horizontal/vertical electric permittivities ϵ_h/ϵ_v (-); $\#epermH = \#epermV = \#res$. Default is ones.

mpermH, mpermV : array_like, optional

Relative horizontal/vertical magnetic permeabilities μ_h/μ_v (-); $\#mpermH = \#mpermV = \#res$. Default is ones.

verb : {0, 1, 2, 3, 4}, optional

Level of verbosity, default is 2:

- 0: Print nothing.
- 1: Print warnings.
- 2: Print additional runtime and kernel calls
- 3: Print additional start/stop, condensed parameter information.
- 4: Print additional full parameter information

Returns PJ0, PJ1 : array

Wavenumber-domain EM responses:

- PJ0: Wavenumber-domain solution for the kernel with a Bessel function of the first kind of order zero.
- PJ1: Wavenumber-domain solution for the kernel with a Bessel function of the first kind of order one.

See also:

dipole Electromagnetic field due to an electromagnetic source (dipoles).

bipole Electromagnetic field due to an electromagnetic source (bipoles).

fem Electromagnetic frequency-domain response.

tem Electromagnetic time-domain response.

Examples

```
>>> import numpy as np
>>> from empymod.model import wavenumber
>>> src = [0, 0, 100]
```



```

>>> rec = [5000, 0, 200]
>>> depth = [0, 300, 1000, 1050]
>>> res = [1e20, .3, 1, 50, 1]
>>> freq = 1
>>> wavenrs = np.logspace(-3.7, -3.6, 10)
>>> PJ0, PJ1 = wavenumber(src, rec, depth, res, freq, wavenrs, verb=0)
>>> print(PJ0)
[ -1.02638329e-08 +4.91531529e-09j -1.05289724e-08 +5.04222413e-09j
  -1.08009148e-08 +5.17238608e-09j -1.10798310e-08 +5.30588284e-09j
  -1.13658957e-08 +5.44279805e-09j -1.16592877e-08 +5.58321732e-09j
  -1.19601897e-08 +5.72722830e-09j -1.22687889e-08 +5.87492067e-09j
  -1.25852765e-08 +6.02638626e-09j -1.29098481e-08 +6.18171904e-09j]
>>> print(PJ1)
[ 1.79483705e-10 -6.59235332e-10j 1.88672497e-10 -6.93749344e-10j
 1.98325814e-10 -7.30068377e-10j 2.08466693e-10 -7.68286748e-10j
 2.19119282e-10 -8.08503709e-10j 2.30308887e-10 -8.50823701e-10j
 2.42062030e-10 -8.95356636e-10j 2.54406501e-10 -9.42218177e-10j
 2.67371420e-10 -9.91530051e-10j 2.80987292e-10 -1.04342036e-09j]

```

`empymod.model.fem(ab, off, angle, zsrc, zrec, lsrc, lrec, depth, freq, etaH, etaV, zetaH, zetaV, xdirect, isfullspace, ht, htarg, use_ne_eval, msrc, mrec, loop_freq, loop_off, conv=True)`
Return the electromagnetic frequency-domain response.

This function is called from one of the above modelling routines. No input-check is carried out here. See the main description of `model` for information regarding input and output parameters.

This function can be directly used if you are sure the provided input is in the correct format. This is useful for inversion routines and similar, as it can speed-up the calculation by omitting input-checks.

`empymod.model.tem(fEM, off, freq, time, signal, ft, ftarg, conv=True)`
Return the time-domain response of the frequency-domain response `fEM`.

This function is called from one of the above modelling routines. No input-check is carried out here. See the main description of `model` for information regarding input and output parameters.

This function can be directly used if you are sure the provided input is in the correct format. This is useful for inversion routines and similar, as it can speed-up the calculation by omitting input-checks.

3.5.2 kernel – Kernel calculation

Kernel of `empymod`, calculates the wavenumber-domain electromagnetic response. Plus analytical full- and half-space solutions.

The functions `wavenumber`, `angle_factor`, `fullspace`, `greenfct`, `reflections`, and `fields` are based on source files (specified in each function) from the source code distributed with [\[Hunziker_et_al_2015\]](#), which can be found at software.seg.org/2015/0001. These functions are (c) 2015 by Hunziker et al. and the Society of Exploration Geophysicists, <http://software.seg.org/disclaimer.txt>. Please read the NOTICE-file in the root directory for more information regarding the involved licenses.

`empymod.kernel.wavenumber(zsrc, zrec, lsrc, lrec, depth, etaH, etaV, zetaH, zetaV, lambd, ab, xdirect, msrc, mrec, use_ne_eval)`
Calculate wavenumber domain solution.

Return the wavenumber domain solutions `PJ0`, `PJ1`, and `PJ0b`, which have to be transformed with a Hankel transform to the frequency domain. `PJ0/PJ0b` and `PJ1` have to be transformed with Bessel functions of order 0 (J_0) and 1 (J_1), respectively.

This function corresponds loosely to equations 105–107, 111–116, 119–121, and 123–128 in [\[Hunziker_et_al_2015\]](#), and equally loosely to the file `kxwmod.c`.

[\[Hunziker_et_al_2015\]](#) uses Bessel functions of orders 0, 1, and 2 (J_0 , J_1 , J_2). The implementations of the *Fast Hankel Transform* and the *Quadrature-with-Extrapolation* in `transform.c` are set-up with Bessel

functions of order 0 and 1 only. This is achieved by applying the recurrence formula

$$J_2(kr) = \frac{2}{kr} J_1(kr) - J_0(kr) .$$

Note: `PJ0` and `PJ0b` could theoretically be added here into one, and then be transformed in one go. However, `PJ0b` has to be multiplied by `factAng` later. This has to be done after the Hankel transform for methods which make use of spline interpolation, in order to work for offsets that are not in line with each other.

This function is called from one of the Hankel functions in `transform`. Consult the modelling routines in `model` for a description of the input and output parameters.

If you are solely interested in the wavenumber-domain solution you can call this function directly. However, you have to make sure all input arguments are correct, as no checks are carried out here.

`empymod.kernel.angle_factor` (*angle, ab, msrc, mrec*)

Return the angle-dependent factor.

The whole calculation in the wavenumber domain is only a function of the distance between the source and the receiver, it is independent of the angel. The angle-dependency is this factor, which can be applied to the corresponding parts in the wavenumber or in the frequency domain.

The `angle_factor` corresponds to the sine and cosine-functions in Eqs 105-107, 111-116, 119-121, 123-128.

This function is called from one of the Hankel functions in `transform`. Consult the modelling routines in `model` for a description of the input and output parameters.

`empymod.kernel.fullspace` (*off, angle, zsrc, zrec, etaH, etaV, zetaH, zetaV, ab, msrc, mrec*)

Analytical full-space solutions in the frequency domain.

$$\hat{G}_{\alpha\beta}^{ee}, \hat{G}_{3\alpha}^{ee}, \hat{G}_{33}^{ee}, \hat{G}_{\alpha\beta}^{em}, \hat{G}_{\alpha 3}^{em}$$

This function corresponds to equations 45–50 in [Hunziker_et_al_2015], and loosely to the corresponding files `Gin11.F90`, `Gin12.F90`, `Gin13.F90`, `Gin22.F90`, `Gin23.F90`, `Gin31.F90`, `Gin32.F90`, `Gin33.F90`, `Gin41.F90`, `Gin42.F90`, `Gin43.F90`, `Gin51.F90`, `Gin52.F90`, `Gin53.F90`, `Gin61.F90`, and `Gin62.F90`.

This function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

`empymod.kernel.greenfct` (*zsrc, zrec, lsrc, lrec, depth, etaH, etaV, zetaH, zetaV, lambda, ab, xdirect, msrc, mrec, use_ne_eval*)

Calculate Green's function for TM and TE.

$$\tilde{g}_{hh}^{tm}, \tilde{g}_{hz}^{tm}, \tilde{g}_{zh}^{tm}, \tilde{g}_{zz}^{tm}, \tilde{g}_{hh}^{te}, \tilde{g}_{zz}^{te}$$

This function corresponds to equations 108–110, 117/118, 122; 89–94, A18–A23, B13–B15; 97–102 A26–A31, and B16–B18 in [Hunziker_et_al_2015], and loosely to the corresponding files `Gamma.F90`, `Wprop.F90`, `Ptotalx.F90`, `Ptotalxm.F90`, `Ptotaly.F90`, `Ptotalym.F90`, `Ptotalz.F90`, and `Ptotalzm.F90`.

The Green's functions are multiplied according to Eqs 105-107, 111-116, 119-121, 123-128; with the factors inside the integrals.

This function is called from the function `kernel.wavenumber`.

`empymod.kernel.reflections` (*depth, e_zH, Gam, lrec, lsrc, use_ne_eval*)

Calculate R_p , R_m .

$$R_n^{\pm}, \bar{R}_n^{\pm}$$

This function corresponds to equations 64/65 and A-11/A-12 in [Hunziker_et_al_2015], and loosely to the corresponding files `Rmin.F90` and `Rplus.F90`.

This function is called from the function `kernel.greenfct`.

`empymod.kernel.fields` (*depth, Rp, Rm, Gam, lrec, lsrc, zsrc, ab, TM, use_ne_eval*)
Calculate Pu+, Pu-, Pd+, Pd-.

$$P_s^{u\pm}, P_s^{d\pm}, \bar{P}_s^{u\pm}, \bar{P}_s^{d\pm}, P_{s-1}^{u\pm}, P_n^{u\pm}, \bar{P}_{s-1}^{u\pm}, \bar{P}_n^{u\pm}, P_{s+1}^{d\pm}, P_n^{d\pm}, \bar{P}_{s+1}^{d\pm}, \bar{P}_n^{d\pm}$$

This function corresponds to equations 81/82, 95/96, 103/104, A-8/A-9, A-24/A-25, and A-32/A-33 in [Hunziker_et_al_2015], and loosely to the corresponding files `Pdownmin.F90`, `Pdownplus.F90`, `Pupmin.F90`, and `Pdownmin.F90`.

This function is called from the function `kernel.greenfct`.

`empymod.kernel.halfspace` (*off, angle, zsrc, zrec, etaH, etaV, freqtime, ab, signal, solution='dhs'*)
Return frequency- or time-space domain VTI half-space solution.

Calculates the frequency- or time-space domain electromagnetic response for a half-space below air using the diffusive approximation, as given in [Slob_et_al_2010], where the electric source is located at [0, 0, zsrc], and the electric receiver at [xco, yco, zrec].

It can also be used to calculate the fullspace solution or the separate fields: direct field, reflected field, and airwave; always using the diffusive approximation. See `solution`-parameter.

This function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and solution parameters.

3.5.3 transform – Hankel and Fourier Transforms

Methods to carry out the required Hankel transform from wavenumber to frequency domain and Fourier transform from frequency to time domain.

The functions for the QWE and DLF Hankel and Fourier transforms are based on source files (specified in each function) from the source code distributed with [Key_2012], which can be found at software.seg.org/2012/0003. These functions are (c) 2012 by Kerry Key and the Society of Exploration Geophysicists, <http://software.seg.org/disclaimer.txt>. Please read the NOTICE-file in the root directory for more information regarding the involved licenses.

`empymod.transform.fht` (*zsrc, zrec, lsrc, lrec, off, angle, depth, ab, etaH, etaV, zetaH, zetaV, xdirect, fhtarg, use_ne_eval, msrc, mrec*)
Hankel Transform using the Digital Linear Filter method.

The *Digital Linear Filter* method was introduced to geophysics by [Ghosh_1970], and made popular and wide-spread by [Anderson_1975], [Anderson_1979], [Anderson_1982]. The DLF is sometimes referred to as the *Fast Hankel Transform* FHT, from which this routine has its name.

This implementation of the DLF follows [Key_2012], equation 6. Without going into the mathematical details (which can be found in any of the above papers) and following [Key_2012], the DLF method rewrites the Hankel transform of the form

$$F(r) = \int_0^\infty f(\lambda) J_v(\lambda r) d\lambda$$

as

$$F(r) = \sum_{i=1}^n f(b_i/r) h_i/r,$$

where h is the digital filter. The Filter abscissae b is given by

$$b_i = \lambda_i r = e^{ai}, \quad i = -l, -l+1, \dots, l,$$

with $l = (n-1)/2$, and a is the spacing coefficient.

This function is loosely based on `get_CSEM1D_FD_FHT.m` from the source code distributed with [\[Key_2012\]](#).

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

Returns fEM : array
Returns frequency-domain EM response.
kcount : int
Kernel count. For DLF, this is 1.
conv : bool
Only relevant for QWE/QUAD.

`empymod.transform.hqwe` (*zsrc, zrec, lsrc, lrec, off, angle, depth, ab, etaH, etaV, zetaH, zetaV, xdirect, qweargs, use_ne_eval, msrc, mrec*)

Hankel Transform using Quadrature-With-Extrapolation.

Quadrature-With-Extrapolation was introduced to geophysics by [\[Key_2012\]](#). It is one of many so-called *ISE* methods to solve Hankel Transforms, where *ISE* stands for Integration, Summation, and Extrapolation.

Following [\[Key_2012\]](#), but without going into the mathematical details here, the QWE method rewrites the Hankel transform of the form

$$F(r) = \int_0^{\infty} f(\lambda) J_v(\lambda r) d\lambda$$

as a quadrature sum which form is similar to the DLF (equation 15),

$$F_i \approx \sum_{j=1}^m f(x_j/r) w_j g(x_j) = \sum_{j=1}^m f(x_j/r) \hat{g}(x_j),$$

but with various bells and whistles applied (using the so-called Shanks transformation in the form of a routine called ϵ -algorithm ([\[Shanks_1955\]](#), [\[Wynn_1956\]](#); implemented with algorithms from [\[Trefethen_2000\]](#) and [\[Weniger_1989\]](#)).

This function is based on `get_CSEM1D_FD_QWE.m`, `qwe.m`, and `getBesselWeights.m` from the source code distributed with [\[Key_2012\]](#).

In the spline-version, `hqwe` checks how steep the decay of the wavenumber-domain result is, and calls QUAD for the very steep interval, for which QWE is not suited.

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

Returns fEM : array
Returns frequency-domain EM response.
kcount : int
Kernel count.
conv : bool
If true, QWE/QUAD converged. If not, <htarg> might have to be adjusted.

`empymod.transform.hquad` (*zsrc, zrec, lsrc, lrec, off, angle, depth, ab, etaH, etaV, zetaH, zetaV, xdirect, quadargs, use_ne_eval, msrc, mrec*)

Hankel Transform using the QUADPACK library.

This routine uses the `scipy.integrate.quad` module, which in turn makes use of the Fortran library QUADPACK (`qagse`).

It is massively (orders of magnitudes) slower than either `fht` or `hqwe`, and is mainly here for completeness and comparison purposes. It always uses interpolation in the wavenumber domain, hence it generally will not be as precise as the other methods. However, it might work in some areas where the others fail.

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

Returns fEM : array
Returns frequency-domain EM response.
kcount : int

Kernel count. For HQUAD, this is 1.

conv : bool

If true, QUAD converged. If not, <htarg> might have to be adjusted.

`empymod.transform.fhft(fEM, time, freq, ftarg)`

Fourier Transform using the Digital Linear Filter method.

It follows the Filter methodology [Anderson_1975], using Cosine- and Sine-filters; see `fhft` for more information.

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

This function is based on `get_CSEM1D_TD_FHT.m` from the source code distributed with [Key_2012].

Returns tEM : array

Returns time-domain EM response of `fEM` for given `time`.

conv : bool

Only relevant for QWE/QUAD.

`empymod.transform.fqwe(fEM, time, freq, qweargs)`

Fourier Transform using Quadrature-With-Extrapolation.

It follows the QWE methodology [Key_2012] for the Hankel transform, see `hqwe` for more information.

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

This function is based on `get_CSEM1D_TD_QWE.m` from the source code distributed with [Key_2012].

`fqwe` checks how steep the decay of the frequency-domain result is, and calls QUAD for the very steep interval, for which QWE is not suited.

Returns tEM : array

Returns time-domain EM response of `fEM` for given `time`.

conv : bool

If true, QWE/QUAD converged. If not, <ftarg> might have to be adjusted.

`empymod.transform.fftlog(fEM, time, freq, ftarg)`

Fourier Transform using FFTLog.

FFTLog is the logarithmic analogue to the Fast Fourier Transform FFT. FFTLog was presented in Appendix B of [Hamilton_2000] and published at <<http://casa.colorado.edu/~ajsh/FFTLog>>.

This function uses a simplified version of `pyfftlog`, which is a python-version of FFTLog. For more details regarding `pyfftlog` see <<https://github.com/prisae/pyfftlog>>.

Not the full flexibility of FFTLog is available here: Only the logarithmic FFT (`fft1` in FFTLog), not the Hankel transform (`fhft` in FFTLog). Furthermore, the following parameters are fixed:

- `kr` = 1 (initial value)
- `kropt` = 1 (silently adjusts `kr`)
- `dir` = 1 (forward)

Furthermore, `q` is restricted to $-1 \leq q \leq 1$.

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

Returns tEM : array

Returns time-domain EM response of `fEM` for given `time`.

conv : bool

Only relevant for QWE/QUAD.

`empymod.transform.fft(fEM, time, freq, ftarg)`

Fourier Transform using the Fast Fourier Transform.

The function is called from one of the modelling routines in `model`. Consult these modelling routines for a description of the input and output parameters.

Returns tEM : array

Returns time-domain EM response of `fEM` for given `time`.

conv : bool

Only relevant for QWE/QUAD.

`empymod.transform.dlf(signal, points, out_pts, filt, pts_per_dec, kind=None, factAng=None, ab=None)`

Digital Linear Filter method.

This is the kernel of the DLF method, used for the Hankel (`fht`) and the Fourier (`ffht`) Transforms. See `fht` for an extensive description.

For the Hankel transform, *signal* contains 3 complex wavenumber-domain signals: (PJ0, PJ1, PJ0b), as returned from *kernel.wavenumber*. The Hankel DLF has two additional, optional parameters: *factAng*, as returned from *kernel.angle_factor*, and *ab*. The PJ0-kernel is the part of the wavenumber-domain calculation which contains a zeroth-order Bessel function and does NOT depend on the angle between source and receiver, only on offset. PJ0b and PJ1 are the parts of the wavenumber-domain calculation which contain a zeroth- and first-order Bessel function, respectively, and can depend on the angle between source and receiver. PJ0, PJ1, or PJ0b can also be None, if they are not used.

For the Fourier transform, *signal* is a complex frequency-domain signal. The Fourier DLF requires one additional parameter, *kind*, which will be ‘cos’ or ‘sin’.

`empymod.transform.qwe(rtol, atol, maxint, inp, intervals, lambd=None, off=None, factAng=None)`

Quadrature-With-Extrapolation.

This is the kernel of the QWE method, used for the Hankel (`hqwe`) and the Fourier (`fqwe`) Transforms. See `hqwe` for an extensive description.

This function is based on `qwe.m` from the source code distributed with [\[Key_2012\]](#).

`empymod.transform.get_spline_values(filt, inp, nr_per_dec=None)`

Return required calculation points.

`empymod.transform.fhti(rmin, rmax, n, q, mu)`

Return parameters required for FFTLog.

3.5.4 filters – Digital Linear Filters

Filters for the *Digital Linear Filter* (DLF) method for the Hankel [\[Ghosh_1970\]](#) and the Fourier [\(\[Anderson_1975\]\)](#) transforms.

To calculate the `dlf.factor` I used

```
np.around(np.average(dlf.base[1:]/dlf.base[:-1]), 15)
```

The filters `kong_61_2007` and `kong_241_2007` from [\[Kong_2007\]](#), and `key_101_2009`, `key_201_2009`, `key_401_2009`, `key_81_CosSin_2009`, `key_241_CosSin_2009`, and `key_601_CosSin_2009` from [\[Key_2009\]](#) are taken from *DIPOLE1D*, [\[Key_2009\]](#), which can be downloaded at <http://marineemlab.ucsd.edu/Projects/Occam/1DCSEM> (1DCSEM). *DIPOLE1D* is distributed under the license GNU GPL version 3 or later. Kerry Key gave his written permission to re-distribute the filters under the Apache License, Version 2.0 (email from Kerry Key to Dieter Werthmüller, 21 November 2016).

The filters `anderson_801_1982` from [\[Anderson_1982\]](#) and `key_51_2012`, `key_101_2012`, `key_201_2012`, `key_101_CosSin_2012`, and `key_201_CosSin_2012`, all from [\[Key_2012\]](#), are taken from the software distributed with [\[Key_2012\]](#) and available at <http://software.seg.org/2012/0003> (SEG-2012-003). These filters are distributed under the SEG license.

The filter `wer_201_2018` was designed with the add-on `fdesign`, see <https://github.com/empymod/article-fdesign>.

class `empymod.filters.DigitalFilter(name, savename=None)`

Simple Class for Digital Linear Filters.

fromfile (*path*='filters')

Load filter values from ascii-files.

Load filter base and filter coefficients from ascii files in the directory *path*; *path* can be a relative or absolute path.

Examples

```
>>> import empymod
>>> # Create an empty filter;
>>> # Name has to be the base of the text files
>>> filt = empymod.filters.DigitalFilter('my-filter')
>>> # Load the ascii-files
>>> filt.fromfile()
>>> # This will load the following three files:
>>> #     ./filters/my-filter_base.txt
>>> #     ./filters/my-filter_j0.txt
>>> #     ./filters/my-filter_j1.txt
>>> # and store them in filt.base, filt.j0, and filt.j1.
```

tofile (*path*='filters')

Save filter values to ascii-files.

Store the filter base and the filter coefficients in separate files in the directory *path*; *path* can be a relative or absolute path.

Examples

```
>>> import empymod
>>> # Load a filter
>>> filt = empymod.filters.wer_201_2018()
>>> # Save it to pure ascii-files
>>> filt.tofile()
>>> # This will save the following three files:
>>> #     ./filters/wer_201_2018_base.txt
>>> #     ./filters/wer_201_2018_j0.txt
>>> #     ./filters/wer_201_2018_j1.txt
```

`empymod.filters.anderson_801_1982()`

Anderson 801 pt Hankel filter, as published in [Anderson_1982].

Taken from file `wa801Hankel.txt` provided with SEG-2012-003.

License: <http://software.seg.org/disclaimer.txt>.

`empymod.filters.key_101_2009()`

Key 101 pt Hankel filter, as published in [Key_2009].

Taken from file `FilterModules.f90` provided with 1DCSEM.

License: Apache License, Version 2.0,.

`empymod.filters.key_101_2012()`

Key 101 pt Hankel filter, as published in [Key_2012].

Taken from file `kk101Hankel.txt` provided with SEG-2012-003.

License: <http://software.seg.org/disclaimer.txt>.

`empymod.filters.key_101_CosSin_2012()`

Key 101 pt CosSin filter, as published in [Key_2012].

Taken from file `kk101CosSin.txt` provided with SEG-2012-003.

License: <http://software.seg.org/disclaimer.txt>.

`empymod.filters.key_201_2009()`
Key 201 pt Hankel filter, as published in [\[Key_2009\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.key_201_2012()`
Key 201 pt Hankel filter, as published in [\[Key_2012\]](#).
Taken from file `kk201Hankel.txt` provided with SEG-2012-003.
License: <http://software.seg.org/disclaimer.txt>.

`empymod.filters.key_201_CosSin_2012()`
Key 201 pt CosSin filter, as published in [\[Key_2012\]](#).
Taken from file `kk201CosSin.txt` provided with SEG-2012-003.
License: <http://software.seg.org/disclaimer.txt>.

`empymod.filters.key_241_CosSin_2009()`
Key 241 pt CosSin filter, as published in [\[Key_2009\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.key_401_2009()`
Key 401 pt Hankel filter, as published in [\[Key_2009\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.key_51_2012()`
Key 51 pt Hankel filter, as published in [\[Key_2012\]](#).
Taken from file `kk51Hankel.txt` provided with SEG-2012-003.
License: <http://software.seg.org/disclaimer.txt>.

`empymod.filters.key_601_CosSin_2009()`
Key 601 pt CosSin filter, as published in [\[Key_2009\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.key_81_CosSin_2009()`
Key 81 pt CosSin filter, as published in [\[Key_2009\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.kong_241_2007()`
Kong 241 pt Hankel filter, as published in [\[Kong_2007\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.kong_61_2007()`
Kong 61 pt Hankel filter, as published in [\[Kong_2007\]](#).
Taken from file `FilterModules.f90` provided with 1DCSEM.
License: [Apache License, Version 2.0](#).

`empymod.filters.wer_201_2018()`
Werthmüller 201 pt Hankel filter, 2018.
Designed with the empymod add-on `fdesign`, see <https://github.com/empymod/article-fdesign>.

License: [Apache License, Version 2.0](#).

3.5.5 `utils` – Utilites

Utilities for `model` such as checking input parameters.

This module consists of four groups of functions:

0. General settings
1. Class `EMArray`
2. Input parameter checks for modelling
3. Internal utilities

class `empymod.utils.EMArray`

Subclassing an `ndarray`: add *amplitude* <amp> and *phase* <pha>.

Parameters `realpart`: array

1. Real part of input, if input is real or complex.
2. Imaginary part of input, if input is pure imaginary.
3. Complex input.

In cases 2 and 3, `imagpart` must be `None`.

imagpart: array, optional

Imaginary part of input. Defaults to `None`.

Examples

```
>>> import numpy as np
>>> from empymod.utils import EMArray
>>> emvalues = EMArray(np.array([1,2,3]), np.array([1, 0, -1]))
>>> print('Amplitude : ', emvalues.amp)
Amplitude : [ 1.41421356  2.          3.16227766]
>>> print('Phase      : ', emvalues.pha)
Phase      : [ 45.          0.        -18.43494882]
```

Attributes

<code>amp</code>	(ndarray) Amplitude of the input data.
<code>pha</code>	(ndarray) Phase of the input data, in degrees, lag-defined (increasing with increasing offset.) To get lead-defined phases, multiply <code>imagpart</code> by -1 before passing through this function.

`empymod.utils.check_time_only` (*time*, *signal*, *verb*)

Check time and signal parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters `time`: array_like

Times *t* (s).

signal: {None, 0, 1, -1}

Source signal:

- None: Frequency-domain response
- -1 : Switch-off time-domain response
- 0 : Impulse time-domain response

- +1 : Switch-on time-domain response

verb : {0, 1, 2, 3, 4}
Level of verbosity.

Returns time : float
Time, checked for size and assured min_time.

`empymod.utils.check_time` (*time, signal, ft, ftarg, verb*)
Check time domain specific input parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters time : array_like
Times *t* (s).

signal : {None, 0, 1, -1}

Source signal:

- None: Frequency-domain response
- -1 : Switch-off time-domain response
- 0 : Impulse time-domain response
- +1 : Switch-on time-domain response

ft : {'sin', 'cos', 'qwe', 'fftlog', 'fft'}
Flag for Fourier transform.

ftarg : str or filter from `empymod.filters` or array_like,
Only used if `signal != None`. Depends on the value for `ft`:

verb : {0, 1, 2, 3, 4}
Level of verbosity.

Returns time : float
Time, checked for size and assured min_time.

freq : float
Frequencies required for given times and ft-settings.

ft, ftarg
Checked if valid and set to defaults if not provided, checked with `signal`.

`empymod.utils.check_model` (*depth, res, aniso, epermH, epermV, mpermH, mpermV, xdirect, verb*)
Check the model: depth and corresponding layer parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters depth : list

Absolute layer interfaces *z* (m); #depth = #res - 1 (excluding +/- infinity).

res : array_like
Horizontal resistivities ρ_h (Ohm.m); #res = #depth + 1.

aniso : array_like
Anisotropies $\lambda = \sqrt{\rho_v/\rho_h}$ (-); #aniso = #res.

epermH, epermV : array_like
Relative horizontal/vertical electric permittivities ϵ_h/ϵ_v (-);
#epermH = #epermV = #res.

mpermH, mpermV : array_like
Relative horizontal/vertical magnetic permeabilities μ_h/μ_v (-);
#mpermH = #mpermV = #res.

xdirect : bool, optional
If True and source and receiver are in the same layer, the direct field is calculated analytically in the frequency domain, if False it is calculated in the wavenumber domain.

verb : {0, 1, 2, 3, 4}
Level of verbosity.

Returns depth : array
Depths of layer interfaces, adds -infty at beginning if not present.

res : array
As input, checked for size.

aniso : array
As input, checked for size. If None, defaults to an array of ones.

epermH, epermV : array_like
As input, checked for size. If None, defaults to an array of ones.

mpermH, mpermV : array_like
As input, checked for size. If None, defaults to an array of ones.

isfullspace : bool
If True, the model is a fullspace (res, aniso, epermH, epermV, mpermH, and mpermV are in all layers the same).

`empymod.utils.check_frequency(freq, res, aniso, epermH, epermV, mpermH, mpermV, verb)`
Calculate frequency-dependent parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters **freq** : array_like
Frequencies f (Hz).

res : array_like
Horizontal resistivities ρ_h (Ohm.m); $\#res = \#depth + 1$.

aniso : array_like
Anisotropies $\lambda = \sqrt{\rho_v/\rho_h}$ (-); $\#aniso = \#res$.

epermH, epermV : array_like
Relative horizontal/vertical electric permittivities ϵ_h/ϵ_v (-);
 $\#epermH = \#epermV = \#res$.

mpermH, mpermV : array_like
Relative horizontal/vertical magnetic permeabilities μ_h/μ_v (-);
 $\#mpermH = \#mpermV = \#res$.

verb : {0, 1, 2, 3, 4}
Level of verbosity.

Returns **freq** : float
Frequency, checked for size and assured `min_freq`.

etaH, etaV : array
Parameters η_H/η_V , same size as provided resistivity.

zetaH, zetaV : array
Parameters ζ_H/ζ_V , same size as provided resistivity.

`empymod.utils.check_hankel(ht, htarg, verb)`
Check Hankel transform parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters **ht** : {'fht', 'qwe', 'quad'}
Flag to choose the Hankel transform.

htarg : str or filter from `empymod.filters` or array_like,
Depends on the value for `ht`.

verb : {0, 1, 2, 3, 4}
Level of verbosity.

Returns **ht, htarg**
Checked if valid and set to defaults if not provided.

`empymod.utils.check_opt(opt, loop, ht, htarg, verb)`
Check optimization parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters **opt** : {None, 'parallel'}
Optimization flag; use `numexpr` or not.

loop : {None, 'freq', 'off'}
Loop flag.

ht : str
Flag to choose the Hankel transform.
htarg : array_like,
Depends on the value for ht.
verb : {0, 1, 2, 3, 4}
Level of verbosity.
Returns **use_ne_eval** : bool
Boolean if to use numexpr.
loop_freq : bool
Boolean if to loop over frequencies.
loop_off : bool
Boolean if to loop over offsets.

`empymod.utils.check_dipole(inp, name, verb)`
Check dipole parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters **inp** : list of floats or arrays
Pole coordinates (m): [pole-x, pole-y, pole-z].
name : str, {'src', 'rec'}
Pole-type.
verb : {0, 1, 2, 3, 4}
Level of verbosity.
Returns **inp** : list
List of pole coordinates [x, y, z].
ninp : int
Number of inp-elements

`empymod.utils.check_bipole(inp, name)`
Check di-/bipole parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters **inp** : list of floats or arrays
Coordinates of inp (m): [dipole-x, dipole-y, dipole-z, azimuth, dip] or.
[bipole-x0, bipole-x1, bipole-y0, bipole-y1, bipole-z0, bipole-z1].
name : str, {'src', 'rec'}
Pole-type.
Returns **inp** : list
As input, checked for type and length.
ninp : int
Number of inp.
ninpz : int
Number of inp depths (ninpz is either 1 or ninp).
isdipole : bool
True if inp is a dipole.

`empymod.utils.check_ab(ab, verb)`
Check source-receiver configuration.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters **ab** : int
Source-receiver configuration.
verb : {0, 1, 2, 3, 4}
Level of verbosity.
Returns **ab_calc** : int
Adjusted source-receiver configuration using reciprocity.
msrc, mrec : bool
If True, src/rec is magnetic; if False, src/rec is electric.

`empymod.utils.check_solution (solution, signal, ab, msrc, mrec)`

Check required solution with parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters `solution` : str

String to define analytical solution.

`signal` : {None, 0, 1, -1}

Source signal:

- None: Frequency-domain response
- -1 : Switch-off time-domain response
- 0 : Impulse time-domain response
- +1 : Switch-on time-domain response

`msrc, mrec` : bool

True if src/rec is magnetic, else False.

`empymod.utils.get_abs (msrc, mrec, srcazm, srcdip, recazm, recdip, verb)`

Get required ab's for given angles.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters `msrc, mrec` : bool

True if src/rec is magnetic, else False.

`srcazm, recazm` : float

Horizontal source/receiver angle (azimuth).

`srcdip, recdip` : float

Vertical source/receiver angle (dip).

`verb` : {0, 1, 2, 3, 4}

Level of verbosity.

Returns `ab_calc` : array of int

ab's to calculate for this bipole.

`empymod.utils.get_geo_fact (ab, srcazm, srcdip, recazm, recdip, msr, mrec)`

Get required geometrical scaling factor for given angles.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters `ab` : int

Source-receiver configuration.

`srcazm, recazm` : float

Horizontal source/receiver angle.

`srcdip, recdip` : float

Vertical source/receiver angle.

Returns `fact` : float

Geometrical scaling factor.

`empymod.utils.get_azm_dip (inp, iz, ninpz, inpts, isdipole, strength, name, verb)`

Get angles, interpolation weights and normalization weights.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters `inp` : list of floats or arrays

Input coordinates (m):

- [x0, x1, y0, y1, z0, z1] (bipole of finite length)
- [x, y, z, azimuth, dip] (dipole, infinitesimal small)

`iz` : int

Index of current di-/bipole depth (-).

`ninpz` : int

Total number of di-/bipole depths (`ninpz` = 1 or `npinz` = `nsr`) (-).

intpts : int

Number of integration points for bipole (-).

isdipole : bool

Boolean if inp is a dipole.

strength : float, optional

Source strength (A):

- If 0, output is normalized to source and receiver of 1 m length, and source strength of 1 A.
- If != 0, output is returned for given source and receiver length, and source strength.

name : str, {'src', 'rec'}

Pole-type.

verb : {0, 1, 2, 3, 4}

Level of verbosity.

Returns tout : list of floats or arrays

Dipole coordinates x, y, and z (m).

azm : float or array of floats

Horizontal angle (azimuth).

dip : float or array of floats

Vertical angle (dip).

g_w : float or array of floats

Factors from Gaussian interpolation.

intpts : int

As input, checked.

inp_w : float or array of floats

Factors from source/receiver length and source strength.

`empymod.utils.get_off_ang(src, rec, nsrc, nrec, verb)`

Get depths, offsets, angles, hence spatial input parameters.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters src, rec : list of floats or arrays

Source/receiver dipole coordinates x, y, and z (m).

nsrc, nrec : int

Number of sources/receivers (-).

verb : {0, 1, 2, 3, 4}

Level of verbosity.

Returns off : array of floats

Offsets

angle : array of floats

Angles

`empymod.utils.get_layer_nr(inp, depth)`

Get number of layer in which inp resides.

Note: If `zinp` is on a layer interface, the layer above the interface is chosen.

This check-function is called from one of the modelling routines in `model`. Consult these modelling routines for a detailed description of the input parameters.

Parameters inp : list of floats or arrays

Dipole coordinates (m)

depth : array

Depths of layer interfaces.

Returns linp : int or array_like of int

Layer number(s) in which inp resides (plural only if bipole).

zinp : float or array

`inp[2]` (depths).

`empymod.utils.printstartfinish(verb, inp=None, kcount=None)`

Print start and finish with time measure and kernel count.

`empymod.utils.conv_warning(conv, targ, name, verb)`

Print error if QWE/QUAD did not converge at least once.

`empymod.utils.set_minimum(min_freq=None, min_time=None, min_off=None, min_res=None, min_angle=None)`

Set minimum values of parameters.

The given parameters are set to its minimum value if they are smaller.

Parameters `min_freq` : float, optional

Minimum frequency [Hz] (default 1e-20 Hz).

`min_time` : float, optional

Minimum time [s] (default 1e-20 s).

`min_off` : float, optional

Minimum offset [m] (default 1e-3 m). Also used to round src- & rec-coordinates.

`min_res` : float, optional

Minimum horizontal and vertical resistivity [Ohm.m] (default 1e-20).

`min_angle` : float, optional

Minimum angle [-] (default 1e-10).

`empymod.utils.get_minimum()`

Return the current minimum values.

Returns `min_vals` : dict

Dictionary of current minimum values with keys

- `min_freq` : float
- `min_time` : float
- `min_off` : float
- `min_res` : float
- `min_angle` : float

For a full description of these options, see `set_minimum`.

`empymod.utils.spline_backwards_hankel(ht, htarg, opt)`

Check opt if deprecated 'spline' is used.

Returns corrected htarg, opt.

3.6 Add-ons

3.6.1 fdesign – Digital Linear Filter (DLF) design

The add-on `fdesign` can be used to design digital linear filters for the Hankel or Fourier transform, or for any linear transform ([*Ghosh_1970*]). For this included or provided theoretical transform pairs can be used. Alternatively, one can use the EM modeller `empymod` to use the responses to an arbitrary 1D model as numerical transform pair.

More information can be found in the following places:

- The article about `fdesign` is in the repo <https://github.com/empymod/article-fdesign>
- Example notebooks to design a filter can be found in the repo <https://github.com/empymod/example-notebooks>

This filter designing tool uses the direct matrix inversion method as described in [*Kong_2007*] and is based on scripts by [*Key_2012*]. The whole project of `fdesign` started with the Matlab scripts from Kerry Key, which he used to design his filters for [*Key_2009*], [*Key_2012*]. Fruitful discussions with Evert Slob and Kerry Key improved the add-on substantially.

Note that the use of `empymod` to create numerical transform pairs is, as of now, only implemented for the Hankel transform.

Implemented analytical transform pairs

The following tables list the transform pairs which are implemented by default. Any other transform pair can be provided as input. A transform pair is defined in the following way:

```
from empymod.scripts import fdesign

def my_tp_pair(var):
    '''My transform pair.'''

    def lhs(l):
        return func(l, var)

    def rhs(r):
        return func(r, var)

    return fdesign.Ghosh(name, lhs, rhs)
```

Here, `name` must be one of `j0`, `j1`, `sin`, or `cos`, depending what type of transform pair it is. Additional variables are provided with `var`. The evaluation points of the `lhs` are denoted by `l`, and the evaluation points of the `rhs` are denoted as `r`. As an example here the implemented transform pair `j0_1`

```
def j0_1(a=1):
    '''Hankel transform pair J0_1 ([Anderson_1975]).'''

    def lhs(l):
        return l*np.exp(-a*l**2)

    def rhs(r):
        return np.exp(-r**2/(4*a))/(2*a)

    return Ghosh('j0', lhs, rhs)
```

Implemented Hankel transforms

- `j0_1` [Anderson_1975]

$$\int_0^{\infty} l \exp(-al^2) J_0(lr) dl = \frac{\exp\left(\frac{-r^2}{4a}\right)}{2a}$$

- `j0_2` [Anderson_1975]

$$\int_0^{\infty} \exp(-al) J_0(lr) dl = \frac{1}{\sqrt{a^2 + r^2}}$$

- `j0_3` [Guptasarma_and_Singh_1997]

$$\int_0^{\infty} l \exp(-al) J_0(lr) dl = \frac{a}{(a^2 + r^2)^{3/2}}$$

- `j0_4` [Chave_and_Cox_1982]

$$\int_0^{\infty} \frac{l}{\beta} \exp(-\beta z_v) J_0(lr) dl = \frac{\exp(-\gamma R)}{R}$$

- j0_5 [Chave_and_Cox_1982]

$$\int_0^{\infty} l \exp(-\beta z_v) J_0(lr) dl = \frac{z_v(\gamma R + 1)}{R^3} \exp(-\gamma R)$$

- j1_1 [Anderson_1975]

$$\int_0^{\infty} l^2 \exp(-al^2) J_1(lr) dl = \frac{r}{4a^2} \exp\left(-\frac{r^2}{4a}\right)$$

- j1_2 [Anderson_1975]

$$\int_0^{\infty} \exp(-al) J_1(lr) dl = \frac{\sqrt{a^2 + r^2} - a}{r\sqrt{a^2 + r^2}}$$

- j1_3 [Anderson_1975]

$$\int_0^{\infty} l \exp(-al) J_1(lr) dl = \frac{r}{(a^2 + r^2)^{3/2}}$$

- j1_4 [Chave_and_Cox_1982]

$$\int_0^{\infty} \frac{l^2}{\beta} \exp(-\beta z_v) J_1(lr) dl = \frac{r(\gamma R + 1)}{R^3} \exp(-\gamma R)$$

- j1_5 [Chave_and_Cox_1982]

$$\int_0^{\infty} l^2 \exp(-\beta z_v) J_1(lr) dl = \frac{r z_v (\gamma^2 R^2 + 3\gamma R + 3)}{R^5} \exp(-\gamma R)$$

Where

$$a > 0, r > 0$$

$$z_v = |z_{rec} - z_{src}|$$

$$R = \sqrt{r^2 + z_v^2}$$

$$\gamma = \sqrt{2j\pi\mu_0 f/\rho}$$

$$\beta = \sqrt{l^2 + \gamma^2}$$

Implemented Fourier transforms

- sin_1 [Anderson_1975]

$$\int_0^{\infty} l \exp(-a^2 l^2) \sin(lr) dl = \frac{\sqrt{\pi} r}{4a^3} \exp\left(-\frac{r^2}{4a^2}\right)$$

- sin_2 [Anderson_1975]

$$\int_0^{\infty} \exp(-al) \sin(lr) dl = \frac{r}{a^2 + r^2}$$

- sin_3 [Anderson_1975]

$$\int_0^{\infty} \frac{l}{a^2 + l^2} \sin(lr) dl = \frac{\pi}{2} \exp(-ar)$$

- `cos_1` [Anderson_1975]

$$\int_0^{\infty} \exp(-a^2 l^2) \cos(lr) dl = \frac{\sqrt{\pi}}{2a} \exp\left(-\frac{r^2}{4a^2}\right)$$

- `cos_2` [Anderson_1975]

$$\int_0^{\infty} \exp(-al) \cos(lr) dl = \frac{a}{a^2 + r^2}$$

- `cos_3` [Anderson_1975]

$$\int_0^{\infty} \frac{1}{a^2 + l^2} \cos(lr) dl = \frac{\pi}{2a} \exp(-ar)$$

`empymod.scripts.fdesign.design`(*n*, *spacing*, *shift*, *fI*, *fC*=*False*, *r*=*None*, *r_def*=(1, 1, 2), *reim*=*None*, *cvar*='amp', *error*=0.01, *name*=*None*, *full_output*=*False*, *finish*=*False*, *save*=*True*, *path*='filters', *verb*=2, *plot*=1)

Digital linear filter (DLF) design

This routine can be used to design digital linear filters for the Hankel or Fourier transform, or for any linear transform ([Ghosh_1970]). For this included or provided theoretical transform pairs can be used. Alternatively, one can use the EM modeller `empymod` to use the responses to an arbitrary 1D model as numerical transform pair.

This filter designing tool uses the direct matrix inversion method as described in [Kong_2007] and is based on scripts by [Key_2012]. The tool is an add-on to the electromagnetic modeller `empymod` [Werthmuller_2017]. Fruitful discussions with Evert Slob and Kerry Key improved the add-on substantially.

Example notebooks of its usage can be found in the repo github.com/empymod/example-notebooks.

Parameters *n* : int

Filter length.

spacing: float or tuple (start, stop, num)

Spacing between filter points. If tuple, it corresponds to the input for `np.linspace` with `endpoint=True`.

shift: float or tuple (start, stop, num)

Shift of base from zero. If tuple, it corresponds to the input for `np.linspace` with `endpoint=True`.

fI, fC : transform pairs

Theoretical or numerical transform pair(s) for the inversion (I) and for the check of goodness (fC). fC is optional. If not provided, fI is used for both fI and fC.

r : array, optional

Right-hand side evaluation points for the check of goodness (fC). Defaults to `r = np.logspace(0, 5, 1000)`, which are a lot of evaluation points, and depending on the transform pair way too long r's.

r_def : tuple (add_left, add_right, factor), optional

Definition of the right-hand side evaluation points *r* of the inversion. *r* is derived from the base values, default is (1, 1, 2).

- `rmin = log10(1/max(base)) - add_left`

- `rmax = log10(1/min(base)) + add_right`

- `r = logspace(rmin, rmax, factor*n)`

reim : np.real or np.imag, optional

Which part of complex transform pairs is used for the inversion. Defaults to `np.real`.

cvar : string {'amp', 'r'}, optional

If 'amp', the inversion minimizes the amplitude. If 'r', the inversion maximizes the right-hand side evaluation point r . Defaults is 'amp'.

error : float, optional

Up to which relative error the transformation is considered good in the evaluation of the goodness. Default is 0.01 (1 %).

name : str, optional

Name of the filter. Defaults to `dlf_+str(n)`.

full_output : bool, optional

If True, returns best filter and output from `scipy.optimize.brute`; else only filter. Default is False.

finish : None, True, or callable, optional

If callable, it is passed through to `scipy.optimize.brute`: minimization function to find minimize best result from brute-force approach. Default is None. You can simply provide True in order to use `scipy.optimize.fmin_powell()`. Set this to None if you are only interested in the actually provided spacing/shift-values.

save : bool, optional

If True, best filter is saved to plain text files in `./filters/`. Can be loaded with `fdesign.load_filter(name)`. If full, the inversion output is stored too. You can add `'.gz'` to `name`, which will then save the full inversion output in a compressed file instead of plain text.

path : string, optional

Absolute or relative path where output will be saved if `save=True`. Default is 'filters'.

verb : {0, 1, 2}, optional

Level of verbosity, default is 2:

- 0: Print nothing.
- 1: Print warnings.
- 2: Print additional time, progress, and result

plot : {0, 1, 2, 3}, optional

Level of plot-verbosity, default is 1:

- 0: Plot nothing.
- 1: Plot brute-force result
- 2: Plot additional theoretical transform pairs, and best inv.
- **3: Plot additional inversion result** (can result in lots of plots depending on spacing and shift) If you are using a notebook, use `%matplotlib notebook` to have all inversion results appear in the same plot.

Returns filter : `empymod.filter.DigitalFilter` instance

Best filter for the input parameters.

full : tuple

Output from `scipy.optimize.brute` with `full_output=True`. (Returned when `full_output` is True.)

`empymod.scripts.fdesign.save_filter(name, filt, full=None, path='filters')`

Save DLF-filter and inversion output to plain text files.

`empymod.scripts.fdesign.load_filter(name, full=False, path='filters')`

Load saved DLF-filter and inversion output from text files.

`empymod.scripts.fdesign.plot_result(filt, full, prntres=True)`

QC the inversion result.

Parameters - `filt`, `full` as returned from `fdesign.design` with `full_output=True`

- If `prntres` is True, it calls `fdesign.print_result` as well.

`empymod.scripts.fdesign.print_result` (*filt, full=None*)

Print best filter information.

Parameters - *filt*, *full* as returned from `fdesign.design` with `full_output=True`

class `empymod.scripts.fdesign.Ghosh` (*name, lhs, rhs*)

Simple Class for Theoretical Transform Pairs.

Named after D. P. Ghosh, honouring his 1970 Ph.D. thesis with which he introduced the digital filter method to geophysics (*[Ghosh_1970]*).

`empymod.scripts.fdesign.j0_1` (*a=1*)

Hankel transform pair J0_1 (*[Anderson_1975]*).

`empymod.scripts.fdesign.j0_2` (*a=1*)

Hankel transform pair J0_2 (*[Anderson_1975]*).

`empymod.scripts.fdesign.j0_3` (*a=1*)

Hankel transform pair J0_3 (*[Guptasarma_and_Singh_1997]*).

`empymod.scripts.fdesign.j0_4` (*f=1, rho=0.3, z=50*)

Hankel transform pair J0_4 (*[Chave_and_Cox_1982]*).

Parameters *f*: float

Frequency (Hz)

rho: float

Resistivity (Ohm.m)

z: float

Vertical distance between source and receiver (m)

`empymod.scripts.fdesign.j0_5` (*f=1, rho=0.3, z=50*)

Hankel transform pair J0_5 (*[Chave_and_Cox_1982]*).

Parameters *f*: float

Frequency (Hz)

rho: float

Resistivity (Ohm.m)

z: float

Vertical distance between source and receiver (m)

`empymod.scripts.fdesign.j1_1` (*a=1*)

Hankel transform pair J1_1 (*[Anderson_1975]*).

`empymod.scripts.fdesign.j1_2` (*a=1*)

Hankel transform pair J1_2 (*[Anderson_1975]*).

`empymod.scripts.fdesign.j1_3` (*a=1*)

Hankel transform pair J1_3 (*[Anderson_1975]*).

`empymod.scripts.fdesign.j1_4` (*f=1, rho=0.3, z=50*)

Hankel transform pair J1_4 (*[Chave_and_Cox_1982]*).

Parameters *f*: float

Frequency (Hz)

rho: float

Resistivity (Ohm.m)

z: float

Vertical distance between source and receiver (m)

`empymod.scripts.fdesign.j1_5` (*f=1, rho=0.3, z=50*)

Hankel transform pair J1_5 (*[Chave_and_Cox_1982]*).

Parameters *f*: float

Frequency (Hz)

rho: float

Resistivity (Ohm.m)

z: float

Vertical distance between source and receiver (m)

```
empymod.scripts.fdesign.sin_1 (a=1)
    Fourier sine transform pair sin_1 ([Anderson_1975]).

empymod.scripts.fdesign.sin_2 (a=1)
    Fourier sine transform pair sin_2 ([Anderson_1975]).

empymod.scripts.fdesign.sin_3 (a=1)
    Fourier sine transform pair sin_3 ([Anderson_1975]).

empymod.scripts.fdesign.cos_1 (a=1)
    Fourier cosine transform pair cos_1 ([Anderson_1975]).

empymod.scripts.fdesign.cos_2 (a=1)
    Fourier cosine transform pair cos_2 ([Anderson_1975]).

empymod.scripts.fdesign.cos_3 (a=1)
    Fourier cosine transform pair cos_3 ([Anderson_1975]).

empymod.scripts.fdesign.empy_hankel (ftype, zsrc, zrec, res, freqtime, depth=[],
                                     aniso=None, epermH=None, epermV=None,
                                     mpermH=None, mpermV=None, htarg=None,
                                     verblhs=0, verbrhs=0)
```

Numerical transform pair with empymod.

All parameters except `ftype`, `verblhs`, and `verbrhs` correspond to the input parameters to `empymod.dipole`. See there for more information.

Note that if `depth=[]`, the analytical full-space solutions will be used (much faster).

Parameters `ftype` : str or list of strings

Either of: {'j0', 'j1', 'j2', ['j0', 'j1']}

- 'j0': Analyze J0-term with `ab=11`, `angle=45°`
- 'j1': Analyze J1-term with `ab=31`, `angle=0°`
- 'j2': Analyze J0- and J1-terms jointly with `ab=12`, `angle=45°`
- ['j0', 'j1']: Same as calling `empy_hankel` twice, once with 'j0' and one with 'j1'; can be provided like this to `fdesign.design`.

verblhs, verbrhs: int

verb-values provided to empymod for lhs and rhs.

Note that `ftype='j2'` only works for fC, not for fI.

3.6.2 `tmtemod` – Calculate up- and down-going TM and TE modes

This add-on for empymod adjusts [Hunziker_et_al_2015] for TM/TE-split. The development was initiated by the development of <https://github.com/empymod/csem-ziolkowski-and-slob> ([Ziolkowski_and_Slob_2018]).

This is a stripped-down version of empymod with a lot of simplifications but an important addition. The modeller empymod returns the total field, hence not distinguishing between TM and TE mode, and even less between up- and down-going fields. The reason behind this is simple: The derivation of [Hunziker_et_al_2015], on which empymod is based, returns the total field. In this derivation each mode (TM and TE) contains non-physical contributions. The non-physical contributions have opposite signs in TM and TE, so they cancel each other out in the total field. However, in order to obtain the correct TM and TE contributions one has to remove these non-physical parts.

This is what this routine does, but only for an x-directed electric source with an x-directed electric receiver, and in the frequency domain (`src` and `rec` in same layer). This version of `dipole` returns the signal separated into TM++, TM+-, TM-+, TM-, TE++, TE+-, TE-+, and TE- as well as the direct field TM and TE contributions. The first superscript denotes the direction in which the field diffuses towards the receiver and the second superscript denotes the direction in which the field diffuses away from the source. For both the plus-sign indicates the field diffuses in the downward direction and the minus-sign indicates the field diffuses in the upward direction. It uses empymod wherever possible. See the corresponding functions in empymod for more explanation and documentation regarding input parameters. There are important limitations:

- `ab == 11` [=> x-directed el. source & el. receivers]
- `signal == None` [=> only frequency domain]
- `xdirect == False` [=> direct field calc. in wavenr-domain]
- `ht == 'fht'`
- `htarg == 'key_201_2012'`
- Options `ft`, `ftarg`, `opt`, and `loop` are not available.
- `lsrc == lrec` [=> `src` & `rec` are assumed in same layer!]
- Model must have more than 1 layer
- Electric permittivity and magnetic permeability are isotropic.
- Only one frequency at once.

Theory

The derivation of [Hunziker_et_al_2015], on which `empymod` is based, returns the total field. Internally it also calculates TM and TE modes, and sums these up. However, the separation into TM and TE mode introduces a singularity at $\kappa = 0$. It has no contribution in the space-frequency domain to the total fields, but it introduces non-physical events in each mode with opposite sign (so they cancel each other out in the total field). In order to obtain the correct TM and TE contributions one has to remove these non-physical parts.

To remove the non-physical part we use the file `tmtemod.py` in this directory. This routine is basically a heavily simplified version of `empymod` with the following limitations outlined above.

So `tmtemod.py` returns the signal separated into TM++, TM+-, TM-, TE++, TE+-, TE-, and TE- as well as the direct field TM and TE contributions. The first superscript denotes the direction in which the field diffuses towards the receiver and the second superscript denotes the direction in which the field diffuses away from the source. For both the plus-sign indicates the field diffuses in the downward direction and the minus-sign indicates the field diffuses in the upward direction. The routine uses `empymod` wherever possible, see the corresponding functions in `empymod` for more explanation and documentation regarding input parameters.

Please note that the notation in [Hunziker_et_al_2015] differs from the notation in [Ziolkowski_and_Slob_2018]. I specify therefore always, which notification applies, either *Hun15* or *Zio18*.

We start with equation (105) in *Hun15*:

$$\begin{aligned}\hat{G}_{xx}^{ee}(\mathbf{x}, \mathbf{x}', \omega) &= \hat{G}_{xx;s}^{ee;i}(\mathbf{x} - \mathbf{x}', \omega) + \frac{1}{8\pi} \int_{\kappa=0}^{\infty} \left(\frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} - \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} \right) J_0(\kappa r) \kappa d\kappa \\ &\quad - \frac{\cos(2\phi)}{8\pi} \int_{\kappa=0}^{\infty} \left(\frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} + \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} \right) J_2(\kappa r) \kappa d\kappa.\end{aligned}$$

Ignoring the incident field, and using $J_2 = \frac{2}{\kappa r} J_1 - J_0$ to avoid J_2 -integrals, we get

$$\begin{aligned}\hat{G}_{xx}^{ee}(\mathbf{x}, \mathbf{x}', \omega) &= \frac{1}{8\pi} \int_{\kappa=0}^{\infty} \left(\frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} - \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} \right) J_0(\kappa r) \kappa d\kappa \\ &\quad + \frac{\cos(2\phi)}{8\pi} \int_{\kappa=0}^{\infty} \left(\frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} + \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} \right) J_0(\kappa r) \kappa d\kappa \\ &\quad - \frac{\cos(2\phi)}{4\pi r} \int_{\kappa=0}^{\infty} \left(\frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} + \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} \right) J_1(\kappa r) d\kappa.\end{aligned}$$

From this the TM- and TE-parts follow as

$$\text{TE} = \frac{\cos(2\phi) - 1}{8\pi} \int_{\kappa=0}^{\infty} \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} J_0(\kappa r) \kappa d\kappa - \frac{\cos(2\phi)}{4\pi r} \int_{\kappa=0}^{\infty} \frac{\zeta_s \tilde{g}_{zz;s}^{te}}{\bar{\Gamma}_s} J_1(\kappa r) d\kappa,$$

$$\text{TM} = \frac{\cos(2\phi) + 1}{8\pi} \int_{\kappa=0}^{\infty} \frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} J_0(\kappa r) \kappa d\kappa - \frac{\cos(2\phi)}{4\pi r} \int_{\kappa=0}^{\infty} \frac{\Gamma_s \tilde{g}_{hh;s}^{tm}}{\eta_s} J_1(\kappa r) d\kappa.$$

Equations (108) and (109) in Hun15 yield the required parameters $\tilde{g}_{hh;s}^{tm}$ and $\tilde{g}_{zz;s}^{te}$,

$$\tilde{g}_{hh;s}^{tm} = P_s^{u-} W_s^u + P_s^{d-} W_s^d,$$

$$\tilde{g}_{zz;s}^{te} = \bar{P}_s^{u+} \bar{W}_s^u + \bar{P}_s^{d+} \bar{W}_s^d.$$

The parameters $P_s^{u\pm}$ and $P_s^{d\pm}$ are given in equations (81) and (82), $\bar{P}_s^{u\pm}$ and $\bar{P}_s^{d\pm}$ in equations (A-8) and (A-9); W_s^u and W_s^d in equation (74) in Hun15. This yields

$$\begin{aligned} \tilde{g}_{zz;s}^{te} &= \frac{\bar{R}_s^+}{\bar{M}_s} \left\{ \exp[-\bar{\Gamma}_s(z_s - z + d^+)] + \bar{R}_s^- \exp[-\bar{\Gamma}_s(z_s - z + d_s + d^-)] \right\} \\ &+ \frac{\bar{R}_s^-}{\bar{M}_s} \left\{ \exp[-\bar{\Gamma}_s(z - z_{s-1} + d^-)] + \bar{R}_s^+ \exp[-\bar{\Gamma}_s(z - z_{s-1} + d_s + d^+)] \right\}, \\ &= \frac{\bar{R}_s^+}{\bar{M}_s} \left\{ \exp[-\bar{\Gamma}_s(2z_s - z - z')] + \bar{R}_s^- \exp[-\bar{\Gamma}_s(z' - z + 2d_s)] \right\} \\ &+ \frac{\bar{R}_s^-}{\bar{M}_s} \left\{ \exp[-\bar{\Gamma}_s(z + z' - 2z_{s-1})] + \bar{R}_s^+ \exp[-\bar{\Gamma}_s(z - z' + 2d_s)] \right\}, \end{aligned}$$

where d^\pm is taken from the text below equation (67). There are four terms in the right-hand side, two in the first line and two in the second line. The first term in the first line is the integrand of TE+-, the second term in the first line corresponds to TE++, the first term in the second line is TE-+, and the second term in the second line is TE--.

If we look at TE+-, we have

$$\tilde{g}_{zz;s}^{te+-} = \frac{\bar{R}_s^+}{\bar{M}_s} \exp[-\bar{\Gamma}_s(2z_s - z - z')],$$

and therefore

$$\begin{aligned} \text{TE}^{+-} &= \frac{\cos(2\phi) - 1}{8\pi} \int_{\kappa=0}^{\infty} \frac{\zeta_s \bar{R}_s^+}{\bar{\Gamma}_s \bar{M}_s} \exp[-\bar{\Gamma}_s(2z_s - z - z')] J_0(\kappa r) \kappa d\kappa \\ &- \frac{\cos(2\phi)}{4\pi r} \int_{\kappa=0}^{\infty} \frac{\zeta_s \bar{R}_s^+}{\bar{\Gamma}_s \bar{M}_s} \exp[-\bar{\Gamma}_s(2z_s - z - z')] J_1(\kappa r) d\kappa. \end{aligned}$$

We can compare this to equation (4.165) in Zio18, with $\hat{I}_x^e = 1$ and slightly re-arranging it to look more alike, we get

$$\begin{aligned} \hat{E}_{xx;H}^{+-} &= \frac{y^2}{4\pi r^2} \int_{\kappa=0}^{\infty} \frac{\zeta_1}{\Gamma_1} \frac{R_{H;1}^-}{M_{H;1}} \exp(-\Gamma_1 h^{+-}) J_0(\kappa r) \kappa d\kappa \\ &+ \frac{x^2 - y^2}{4\pi r^3} \int_{\kappa=0}^{\infty} \frac{\zeta_1}{\Gamma_1} \left(\frac{R_{H;1}^-}{M_{H;1}} - \frac{R_{H;1}^-(\kappa=0)}{M_{H;1}(\kappa=0)} \right) \exp(-\Gamma_1 h^{+-}) J_1(\kappa r) d\kappa \\ &- \frac{\zeta_1(x^2 - y^2)}{4\pi \gamma_1 r^4} \frac{R_{H;1}^-(\kappa=0)}{M_{H;1}(\kappa=0)} \exp(-\gamma_1 R^{+-}). \end{aligned}$$

The notation in this equation follows Zio18.

The difference between the two previous equations is that the first one contains non-physical contributions. These have opposite signs in TM+- and TE+-, and therefore cancel each other out. But if we want to know the specific contributions from TM and TE we have to remove them. The non-physical contributions only affect the J_1 -integrals, and only for $\kappa = 0$.

The following lists for all 8 cases the term that has to be removed, in the notation of Zio18 (for the notation as in Hun15 see the implementation in `tmtemod.py`):

$$TE^{++} = + \frac{\zeta_1(x^2 - y^2)}{4\pi \gamma_1 r^4} \frac{\exp(-\gamma_1 |h^-|)}{M_{H;1}(\kappa=0)},$$

$$\begin{aligned}
TE^{-+} &= -\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{R_{H;1}^+(\kappa=0) \exp(-\gamma_1 h^{-+})}{M_{H;1}(\kappa=0)}, \\
TE^{+-} &= -\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{R_{H;1}^-(\kappa=0) \exp(-\gamma_1 h^{+-})}{M_{H;1}(\kappa=0)}, \\
TE^{--} &= +\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{R_{H;1}^+(\kappa=0) R_{H;1}^-(\kappa=0) \exp(-\gamma_1 h^{--})}{M_{H;1}(\kappa=0)}, \\
TM^{++} &= -\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{\exp(-\gamma_1 |h^-|)}{M_{V;1}(\kappa=0)}, \\
TM^{-+} &= -\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{R_{V;1}^+(\kappa=0) \exp(-\gamma_1 h^{-+})}{M_{V;1}(\kappa=0)}, \\
TM^{+-} &= -\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{R_{V;1}^-(\kappa=0) \exp(-\gamma_1 h^{+-})}{M_{V;1}(\kappa=0)}, \\
TM^{--} &= -\frac{\zeta_1(x^2 - y^2)}{4\pi\gamma_1 r^4} \frac{R_{V;1}^+(\kappa=0) R_{V;1}^-(\kappa=0) \exp(-\gamma_1 h^{--})}{M_{V;1}(\kappa=0)}.
\end{aligned}$$

Note that in the first and fourth equations the correction terms have opposite sign as those in the fifth and eighth equations because at $\kappa = 0$ the TM and TE mode correction terms are equal. Also note that in the second and third equations the correction terms have the same sign as those in the sixth and seventh equations because at $\kappa = 0$ the TM and TE mode reflection responses in those terms are equal but with opposite sign: $R_{V;1}^\pm(\kappa=0) = -R_{V;1}^\pm(\kappa=0)$.

Hun15 uses ϕ , whereas Zio18 uses x, y , for which we can use

$$\cos(2\phi) = -\frac{x^2 - y^2}{r^2}.$$

`empymod.scripts.tmtmod.dipole(src, rec, depth, res, freqtime, aniso=None, eperm=None, mperm=None, verb=2)`

Return the electromagnetic field due to a dipole source.

This is a modified version of `empymod.model.dipole()`. It returns the separated contributions of TM-, TM+, TM+, TM++, TMdirect, TE-, TE+, TE+, TE++, and TEDirect.

Parameters `src, rec` : list of floats or arrays

Source and receiver coordinates (m): [x, y, z]. The x- and y-coordinates can be arrays, z is a single value. The x- and y-coordinates must have the same dimension.

Sources or receivers placed on a layer interface are considered in the upper layer.

Sources and receivers must be in the same layer.

depth : list

Absolute layer interfaces z (m); #depth = #res - 1 (excluding +/- infinity).

res : array_like

Horizontal resistivities rho_h (Ohm.m); #res = #depth + 1.

freqtime : float

Frequency f (Hz). (The name `freqtime` is kept for consistency with `empymod.model.dipole()`. Only one frequency at once.

aniso : array_like, optional

Anisotropies $\lambda = \sqrt{\rho_v/\rho_h}$ (-); #aniso = #res. Defaults to ones.

eperm : array_like, optional

Relative electric permittivities epsilon (-); #eperm = #res. Default is ones.

mperm : array_like, optional

Relative magnetic permeabilities mu (-); #mperm = #res. Default is ones.

verb : {0, 1, 2, 3, 4}, optional

Level of verbosity, default is 2:

- 0: Print nothing.
- 1: Print warnings.
- 2: Print additional runtime and kernel calls
- 3: Print additional start/stop, condensed parameter information.
- 4: Print additional full parameter information

Returns **TM, TE** : list of ndarrays, (nfreq, nrec, nsrc)

Frequency-domain EM field [V/m], separated into $TM = [TM-, TM+, TM+-, TM++, TMdirect]$ and $TE = [TE-, TE+, TE+-, TE++, TEdirect]$.

However, source and receiver are normalised. So the source strength is 1 A and its length is 1 m. Therefore the electric field could also be written as $[V/(A.m^2)]$.

The shape of EM is (nfreq, nrec, nsrc). However, single dimensions are removed.

3.6.3 printinfo – Tools to print date, time, and version information

Print or return date, time, and package version information in any environment (Jupyter notebook, IPython console, Python console, QT console), either as html-table (notebook) or as plain text (anywhere).

This script was heavily inspired by

- `ipynbtools.py` from <https://github.com/qutip/qutip>, and
- `watermark.py` from <https://github.com/rasbt/watermark>,

Always shown are the OS, number of CPU(s), numpy, scipy, empymod, sys.version, and time/date.

Additionally shown are, if they can be imported, IPython, matplotlib, and numexpr. It also shows MKL information, if available.

All modules provided in `add_pkg` are also shown. They have to be imported before `versions` is called.

`empymod.scripts.printinfo.versions(mode='print', add_pkg=None, ncol=4)`

Return date, time, and version information.

Print or return date, time, and package version information in any environment (Jupyter notebook, IPython console, Python console, QT console), either as html-table (notebook) or as plain text (anywhere).

This script was heavily inspired by:

- `ipynbtools.py` from `qutip` <https://github.com/qutip/qutip>
- `watermark.py` from <https://github.com/rasbt/watermark>

This is a wrapper for `versions_html` and `versions_text`.

Parameters **mode** : string, optional; {<'print'>, 'HTML', 'Pretty', 'plain', 'html'}

Defaults to 'print':

- 'print': Prints text-version to stdout, nothing returned.
- 'HTML': Returns html-version as `IPython.display.HTML(html)`.
- 'html': Returns html-version as plain text.
- 'Pretty': Returns text-version as `IPython.display.Pretty(text)`.
- 'plain': Returns text-version as plain text.

'HTML' and 'Pretty' require IPython.

add_pkg : packages, optional

Package or list of packages to add to output information (must be imported beforehand).

ncol : int, optional

Number of package-columns in html table; only has effect if mode='HTML' or mode='html'. Defaults to 3.

Returns Depending on mode (HTML-instance; plain text; html as plain text; or nothing, only printing to stdout).

Examples

```
>>> import pytest
>>> import dateutil
>>> from empymod import versions
>>> versions() # Default values
>>> versions('plain', pytest) # Provide additional package
>>> versions('HTML', [pytest, dateutil], ncol=5) # HTML
```

`empymod.scripts.printinfo.versions_html(add_pckg=None, ncol=4)`
HTML version.

See versions for details.

`empymod.scripts.printinfo.versions_text(add_pckg=None)`
Plain-text version.

See versions for details.

Bibliography

- [Anderson_1975] Anderson, W. L., 1975, Improved digital filters for evaluating Fourier and Hankel transform integrals: USGS, PB242800; pubs.er.usgs.gov/publication/70045426.
- [Anderson_1979] Anderson, W. L., 1979, Numerical integration of related Hankel transforms of orders 0 and 1 by adaptive digital filtering: *Geophysics*, 44, 1287–1305; DOI: [10.1190/1.1441007](https://doi.org/10.1190/1.1441007).
- [Anderson_1982] Anderson, W. L., 1982, Fast Hankel transforms using related and lagged convolutions: *ACM Trans. on Math. Softw. (TOMS)*, 8, 344–368; DOI: [10.1145/356012.356014](https://doi.org/10.1145/356012.356014).
- [Chave_and_Cox_1982] Chave, A. D., and C. S. Cox, 1982, Controlled electromagnetic sources for measuring electrical conductivity beneath the oceans: 1. forward problem and model study: *Journal of Geophysical Research*, 87, 5327–5338; DOI: [10.1029/JB087iB07p05327](https://doi.org/10.1029/JB087iB07p05327).
- [Ghosh_1970] Ghosh, D. P., 1970, The application of linear filter theory to the direct interpretation of geoelectrical resistivity measurements: Ph.D. Thesis, TU Delft; UUID: [88a568bb-ebec-4d7b-92df-6639b42da2b2](https://nbn-resolving.org/urn:nbn:nl:uid:88a568bb-ebec-4d7b-92df-6639b42da2b2).
- [Guptasarma_and_Singh_1997] Guptasarma, D., and B. Singh, 1997, New digital linear filters for Hankel J0 and J1 transforms: *Geophysical Prospecting*, 45, 745–762; DOI: [10.1046/j.1365-2478.1997.500292.x](https://doi.org/10.1046/j.1365-2478.1997.500292.x).
- [Haines_and_Jones_1988] Haines, G. V., and A. G. Jones, 1988, Logarithmic Fourier transformation: *Geophysical Journal*, 92, 171–178; DOI: [10.1111/j.1365-246X.1988.tb01131.x](https://doi.org/10.1111/j.1365-246X.1988.tb01131.x).
- [Hamilton_2000] Hamilton, A. J. S., 2000, Uncorrelated modes of the non-linear power spectrum: *Monthly Notices of the Royal Astronomical Society*, 312, pages 257–284; DOI: [10.1046/j.1365-8711.2000.03071.x](https://doi.org/10.1046/j.1365-8711.2000.03071.x); Website of FFTLog: casa.colorado.edu/~ajsh/FFTLog.
- [Hunziker_et_al_2015] Hunziker, J., J. Thorbecke, and E. Slob, 2015, The electromagnetic response in a layered vertical transverse isotropic medium: A new look at an old problem: *Geophysics*, 80(1), F1–F18; DOI: [10.1190/geo2013-0411.1](https://doi.org/10.1190/geo2013-0411.1); Software: software.seg.org/2015/0001.
- [Key_2009] Key, K., 2009, 1D inversion of multicomponent, multifrequency marine CSEM data: Methodology and synthetic studies for resolving thin resistive layers: *Geophysics*, 74(2), F9–F20; DOI: [10.1190/1.3058434](https://doi.org/10.1190/1.3058434). Software: marineemlab.ucsd.edu/Projects/Occam/1DCSEM.
- [Key_2012] Key, K., 2012, Is the fast Hankel transform faster than quadrature?: *Geophysics*, 77(3), F21–F30; DOI: [10.1190/geo2011-0237.1](https://doi.org/10.1190/geo2011-0237.1); Software: software.seg.org/2012/0003.
- [Kong_2007] Kong, F. N., 2007, Hankel transform filters for dipole antenna radiation in a conductive medium: *Geophysical Prospecting*, 55, 83–89; DOI: [10.1111/j.1365-2478.2006.00585.x](https://doi.org/10.1111/j.1365-2478.2006.00585.x).
- [Shanks_1955] Shanks, D., 1955, Non-linear transformations of divergent and slowly convergent sequences: *Journal of Mathematics and Physics*, 34, 1–42; DOI: [10.1002/sapm19553411](https://doi.org/10.1002/sapm19553411).
- [Slob_et_al_2010] Slob, E., J. Hunziker, and W. A. Mulder, 2010, Green’s tensors for the diffusive electric field in a VTI half-space: *PIER*, 107, 1–20; DOI: [10.2528/PIER10052807](https://doi.org/10.2528/PIER10052807).

- [Talman_1978] Talman, J. D., 1978, Numerical Fourier and Bessel transforms in logarithmic variables: *Journal of Computational Physics*, 29, pages 35–48; DOI: [10.1016/0021-9991\(78\)90107-9](https://doi.org/10.1016/0021-9991(78)90107-9).
- [Trefethen_2000] Trefethen, L. N., 2000, Spectral methods in MATLAB: Society for Industrial and Applied Mathematics (SIAM), volume 10 of *Software, Environments, and Tools*, chapter 12, page 129; DOI: [10.1137/1.9780898719598.ch12](https://doi.org/10.1137/1.9780898719598.ch12).
- [Weniger_1989] Weniger, E. J., 1989, Nonlinear sequence transformations for the acceleration of convergence and the summation of divergent series: *Computer Physics Reports*, 10, 189–371; arXiv: [abs/math/0306302](https://arxiv.org/abs/math/0306302).
- [Werthmuller_2017] Werthmüller, D., 2017, An open-source full 3D electromagnetic modeler for 1D VTI media in Python: *empymod: Geophysics*, 82(6), WB9–WB19; DOI: [10.1190/geo2016-0626.1](https://doi.org/10.1190/geo2016-0626.1).
- [Werthmuller_2017b] Werthmüller, D., 2017, Getting started with controlled-source electromagnetic 1D modeling: *The Leading Edge*, 36, 352–355; DOI: [10.1190/tle36040352.1](https://doi.org/10.1190/tle36040352.1).
- [Wynn_1956] Wynn, P., 1956, On a device for computing the $e_m(S_n)$ transformation: *Math. Comput.*, 10, 91–96; DOI: [10.1090/S0025-5718-1956-0084056-6](https://doi.org/10.1090/S0025-5718-1956-0084056-6).
- [Ziolkowski_and_Slob_2018] Ziolkowski, A., and E. Slob, 2018, *Introduction to Controlled-Source Electromagnetic Methods*: Cambridge University Press; expected to be published late 2018.

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